

Piecewise-linear maps: where do they come from and what do they do?

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 - piecewise-linear
 - with two pieces
 - continuous

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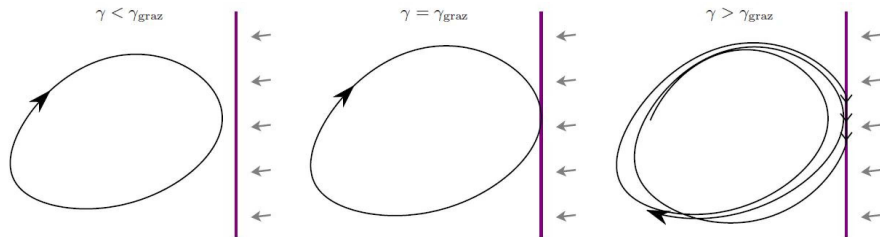
$$\begin{bmatrix} x_1 \\ x_2 \end{bmatrix} \mapsto \begin{bmatrix} -a|x_1| + x_2 + 1 \\ bx_1 \end{bmatrix}$$

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 - with two pieces
 - continuous
- E.g. the Lozi family:

$$\begin{aligned}
 \begin{bmatrix} x_1 \\ x_2 \end{bmatrix} &\mapsto \begin{bmatrix} -a|x_1| + x_2 + 1 \\ bx_1 \end{bmatrix} \\
 &= \begin{cases} \begin{bmatrix} a & 1 \\ b & 0 \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \end{bmatrix} + \begin{bmatrix} 1 \\ 0 \end{bmatrix}, & x_1 \leq 0 \\ \begin{bmatrix} -a & 1 \\ b & 0 \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \end{bmatrix} + \begin{bmatrix} 1 \\ 0 \end{bmatrix}, & x_1 \geq 0 \end{cases}
 \end{aligned}$$

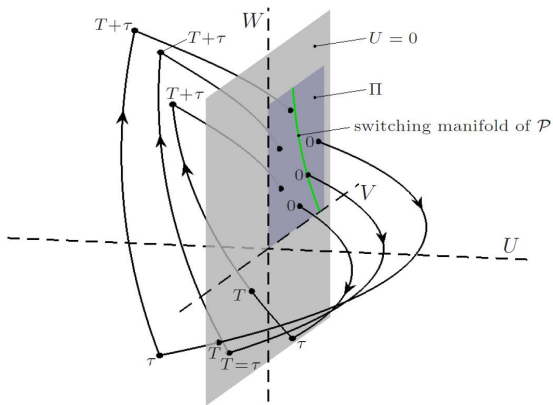
PIECEWISE-SMOOTH MAPS AS POINCARÉ MAPS

1) Grazing-sliding bifurcations



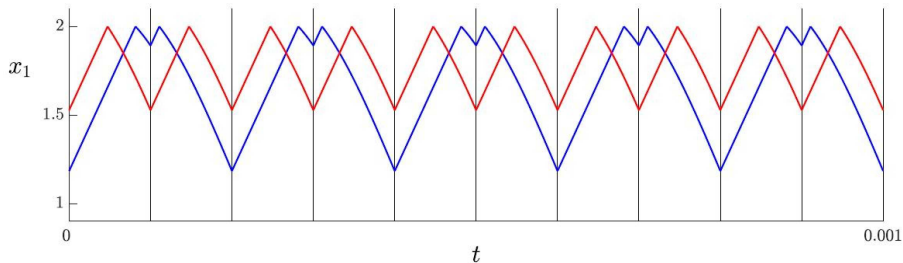
- Get a zero eigenvalue in one piece of the map due to the sliding motion
 - di Bernardo, Kowalczyk & Nordmark, 2002

2) Event collisions for ODEs with time-delayed switching



- Seiber, 2006

3) Corner collisions



- Often get a non-invertible map due to the switching
 - di Bernardo, Budd & Champneys, 2001

BORDER-COLLISION BIFURCATIONS

- Consider a continuous piecewise-smooth map

$$x \mapsto \begin{cases} f_L(x; \mu), & h(x) \leq 0, \\ f_R(x; \mu), & h(x) \geq 0, \end{cases}$$

where $\mu \in \mathbb{R}$ is a parameter.

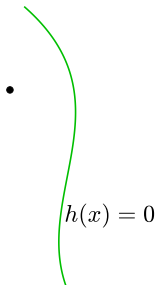

$$h(x) = 0$$

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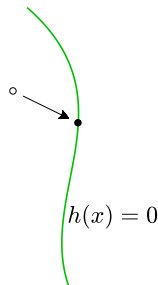
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- As μ is varied a **border-collision bifurcation** occurs when a fixed point collides with the switching manifold.



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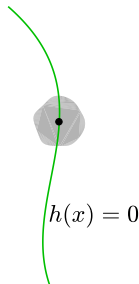
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- As μ is varied a **border-collision bifurcation** occurs when a fixed point collides with the switching manifold.
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$$x \mapsto \begin{cases} A_L x + b\mu, & x_1 \leq 0, \\ A_R x + b\mu, & x_1 \geq 0, \end{cases}$$

where A_L and A_R differ only in their first columns.



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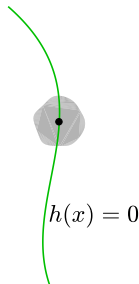
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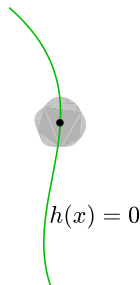
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- Often the truncation to piecewise-linear is justified.



MANY DIFFERENT DYNAMICAL TRANSITIONS

Physica D 57 (1992) 39–57
North-Holland

PHYSICA 

Border-collision bifurcations including “period two to period three” for piecewise smooth systems*

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^b*Vakgroep Econometrie, F.E.W., Rijksuniversiteit Groningen, Postbus 800, 9700 AV Groningen, The Netherlands*

^c*Department of Mathematics, University of Maryland, College Park, MD 20742, USA*

THE BORDER-COLLISION NORMAL FORM

- If (A_L, e_1) is **observable**, equivalently if A_L has no eigenvector v with $v_1 = 0$, then

$$x \mapsto \begin{cases} A_L x + b\mu, & x_1 \leq 0, \\ A_R x + b\mu, & x_1 \geq 0, \end{cases}$$

can be transformed to

$$x \mapsto \begin{cases} C_L x + e_1\mu, & x_1 \leq 0, \\ C_R x + e_1\mu, & x_1 \geq 0, \end{cases}$$

with **companion matrices**:

$$C_L = \begin{bmatrix} p_1 & 1 & & \\ & p_2 & \ddots & \\ & \vdots & & \\ & p_n & & 1 \end{bmatrix}, \quad C_R = \begin{bmatrix} q_1 & 1 & & \\ & q_2 & \ddots & \\ & \vdots & & \\ & q_n & & 1 \end{bmatrix}, \quad e_1 = \begin{bmatrix} 1 \\ 0 \\ \vdots \\ 0 \end{bmatrix}.$$

THE BORDER-COLLISION NORMAL FORM

- In two dimensions

$$x \mapsto \begin{cases} \begin{bmatrix} \tau_L & 1 \\ -\delta_L & 0 \end{bmatrix} x + \begin{bmatrix} 1 \\ 0 \end{bmatrix} \mu, & x_1 \leq 0, \\ \begin{bmatrix} \tau_R & 1 \\ -\delta_R & 0 \end{bmatrix} x + \begin{bmatrix} 1 \\ 0 \end{bmatrix} \mu, & x_1 \geq 0. \end{cases}$$

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- If $\tau_L = -\tau_R = a$, $\delta_L = \delta_R = -b$, and $\mu = 1$, get the Lozi family

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OUTLINE

- 1) A sausage-string structure for periodicity regions.
- 2) Homoclinic corners and infinitely many attractors.
- 3) Stability of boundary fixed points.

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PERIODICITY REGIONS

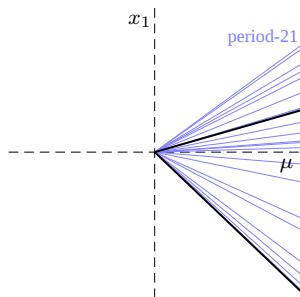
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$$A_L = \begin{bmatrix} 0 & 1 & 0 \\ 1 & 0 & 1 \\ 0.15 & 0 & 0 \end{bmatrix}, \quad A_R = \begin{bmatrix} -1.5 & 1 & 0 \\ 0 & 0 & 1 \\ 2 & 0 & 0 \end{bmatrix}, \quad b = \begin{bmatrix} 1 \\ 0 \\ 0 \end{bmatrix}$$

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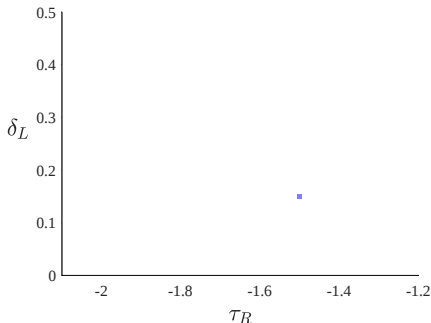
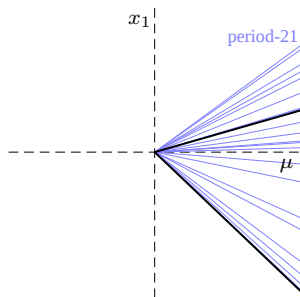


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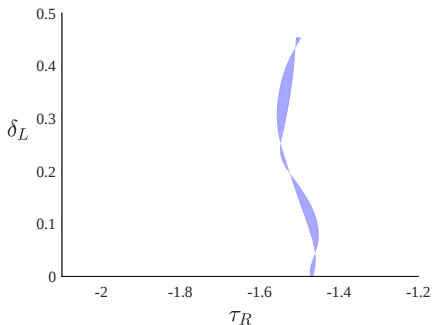
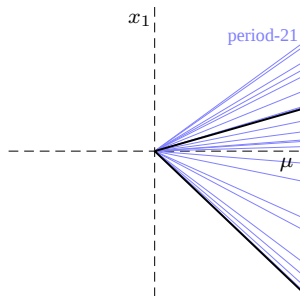


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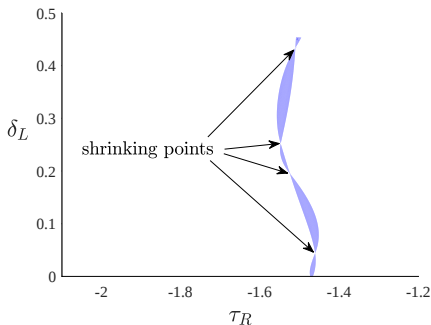
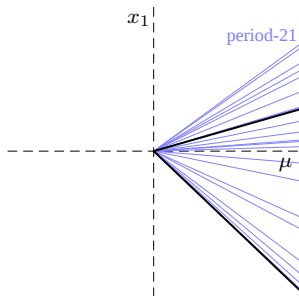


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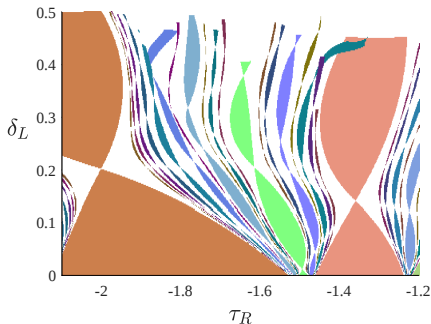
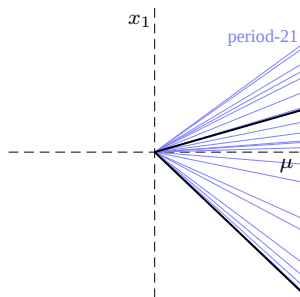


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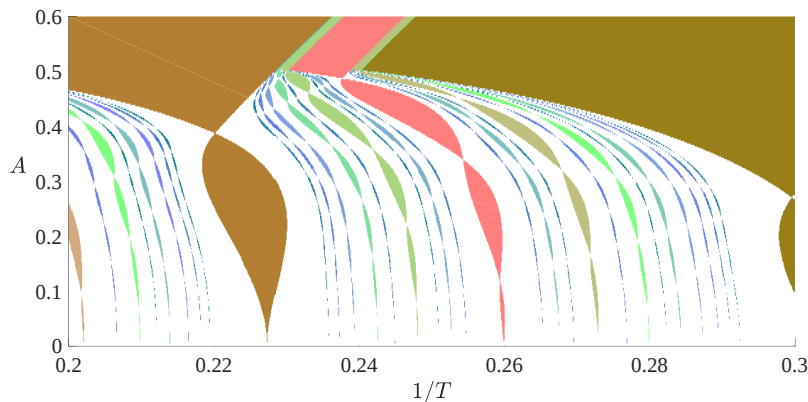


INTEGRATE-AND-FIRE NEURONS

potential: $\dot{V} = -V + I - F(t) \quad (I = 1.5)$

reset law: $V(t) = 1 \rightarrow V(t) = 0$

square-wave forcing: $F(t) = A \operatorname{sgn}[\sin(\frac{2\pi t}{T})]$

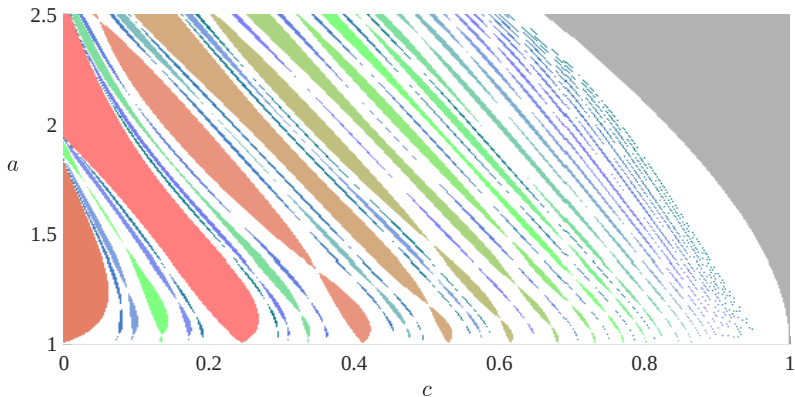


- Tiesinga, 2002

BUSINESS CYCLE DYNAMICS

income: $Y_t = cY_{t-1} + I_t$

net investment: $I_t = \max[a(Y_{t-1} - Y_{t-2}), -arY_{t-2}]$



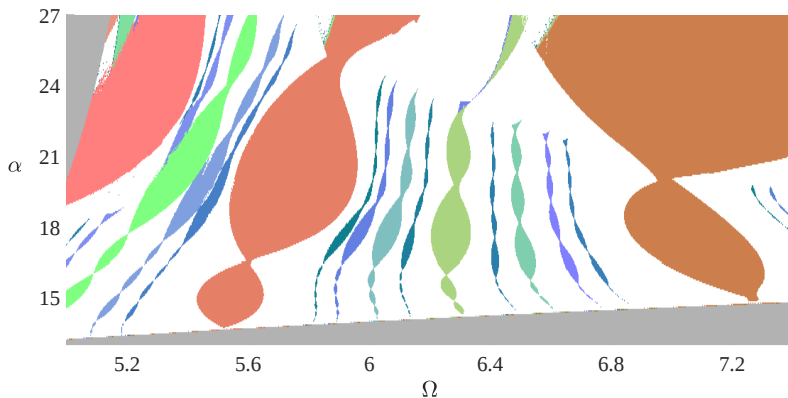
- Sushko, Gardini & Puu, 2004

DC/DC POWER CONVERTERS

$$\dot{y}_k = \lambda_k(y_k - K_F(t)), \quad k = 1, 2$$

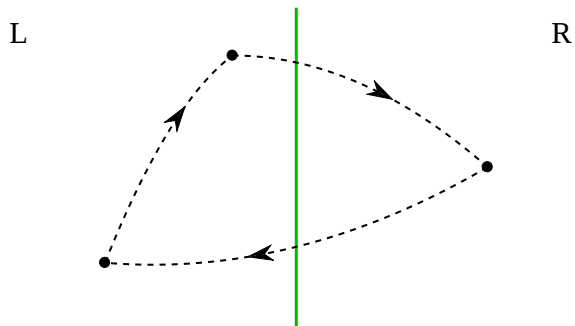
$$K_F(t) = \mathbf{1}_{\xi > \frac{q}{\alpha\Omega}(t - \lfloor t \rfloor)}$$

$$\xi = y_1(\lfloor t \rfloor) - \theta y_2(\lfloor t \rfloor) + \frac{q}{2\Omega}$$

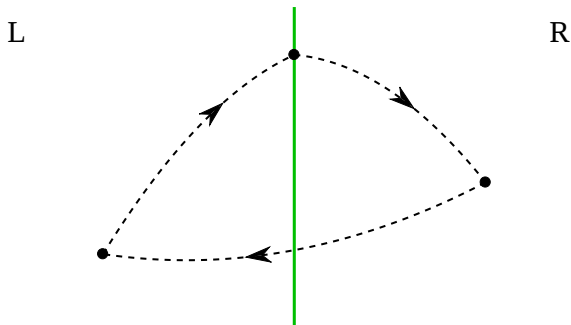


- Zhusubaliyev & Mosekilde, 2008

- An LLR -cycle:



- The left and right boundaries of the periodicity regions are border-collision bifurcations where one point of an $\mathcal{S}$ -cycle lies on the switching manifold.



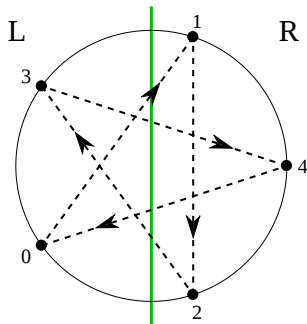
- Shrinking points are where an $\mathcal{S}$ -cycle has two points on the switching manifold, where $\mathcal{S} = \mathcal{F}[\ell, m, n]$ corresponds to rigid rotation on a circle:

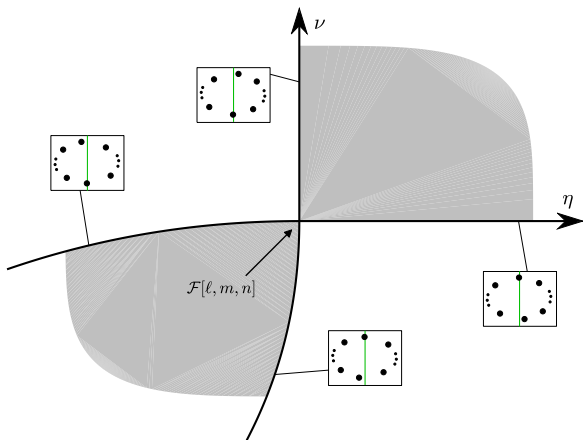
n – period

$\frac{m}{n}$ – rotation number

ℓ – number of L's

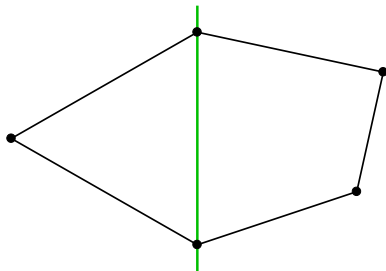
- For example, $\mathcal{F}[2, 2, 5] = LRRLR \rightarrow$





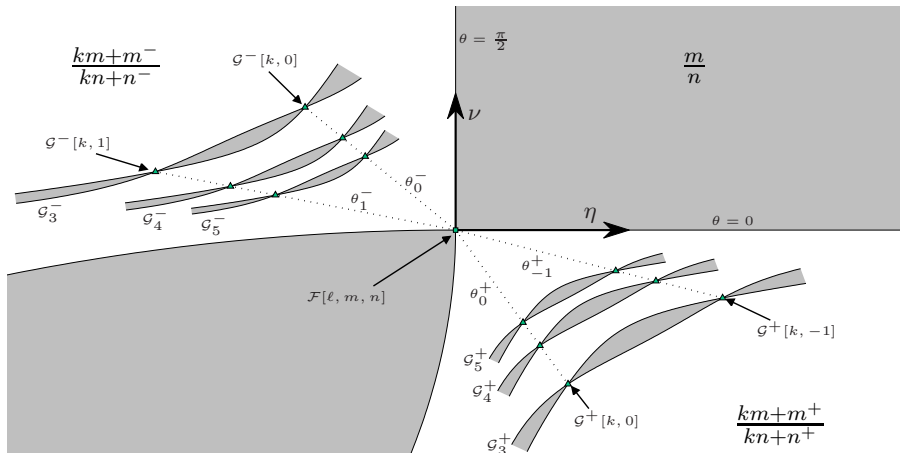
- Four distinct curves of border-collision bifurcations.
- These admit a nice combinatorial characterisation in terms of ℓ , m , and n .
 - S & Meiss, *Nonlinearity*, 2009 & 2010

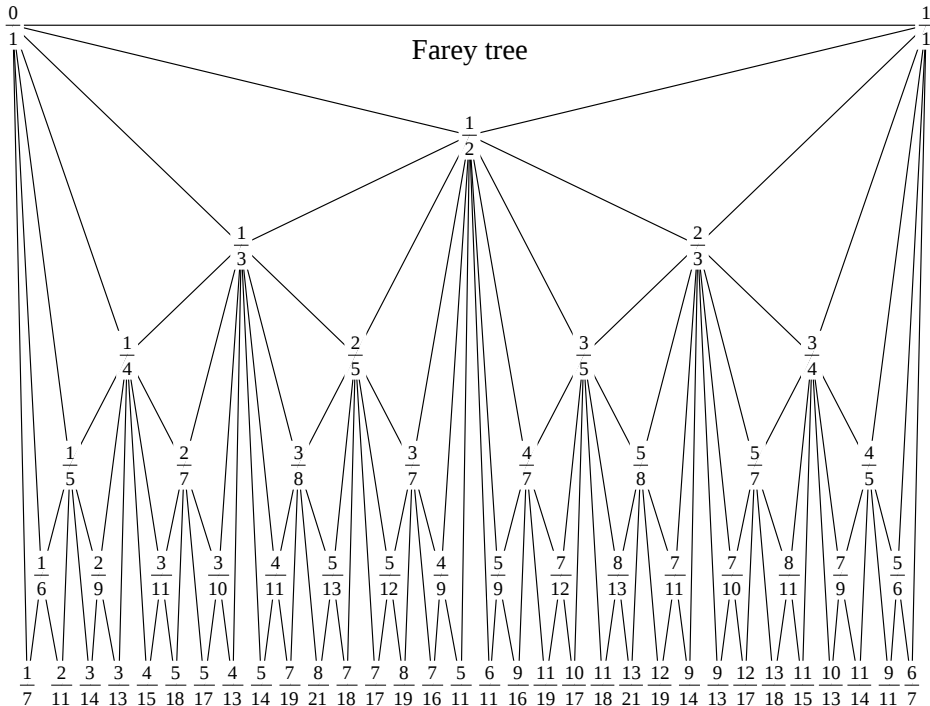
- At a shrinking point there exists an invariant polygon.

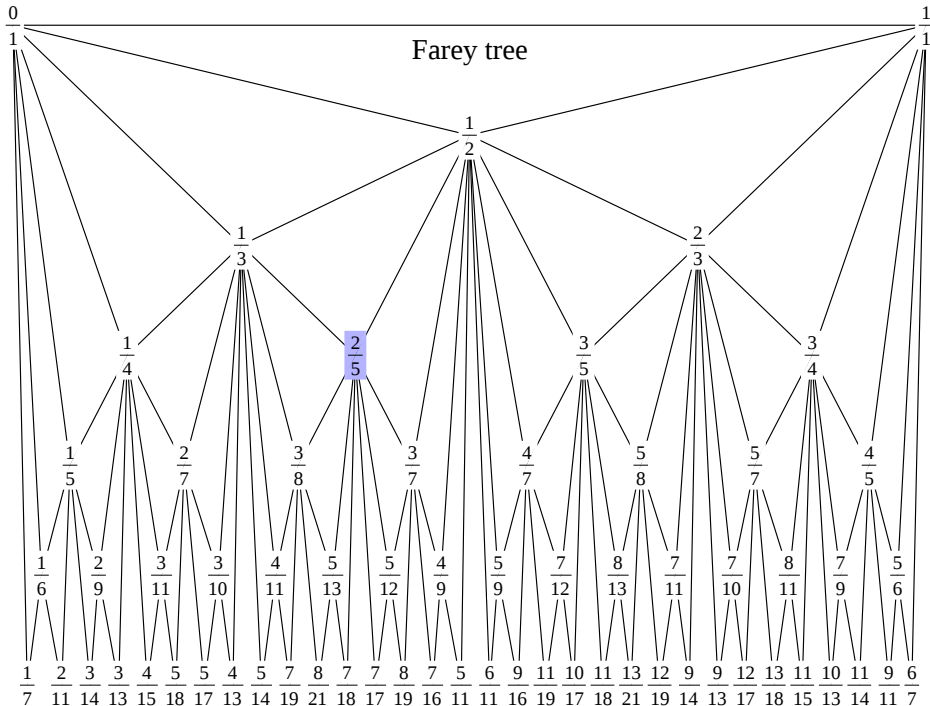


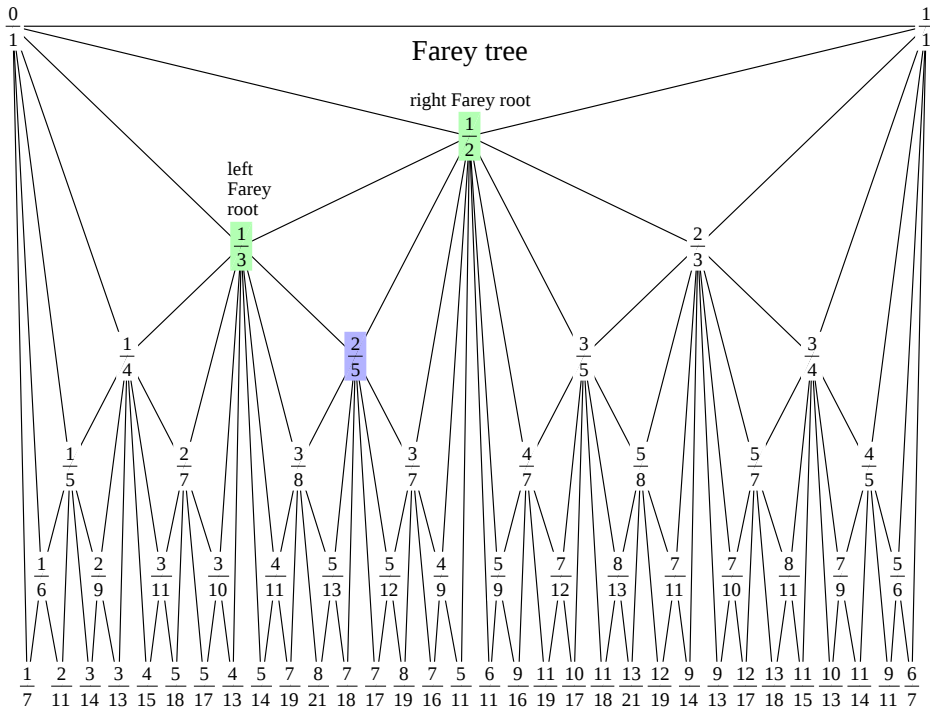
- Near the shrinking point this persists as a one-dimensional slow manifold.

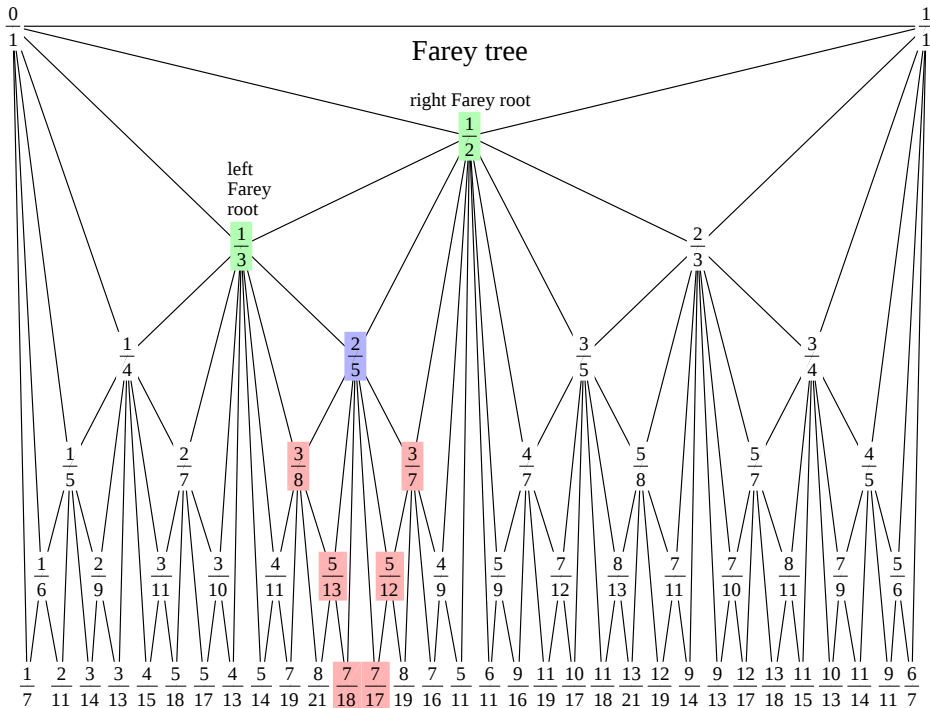
- Near a shrinking point there are two primary sequences of periodicity regions.



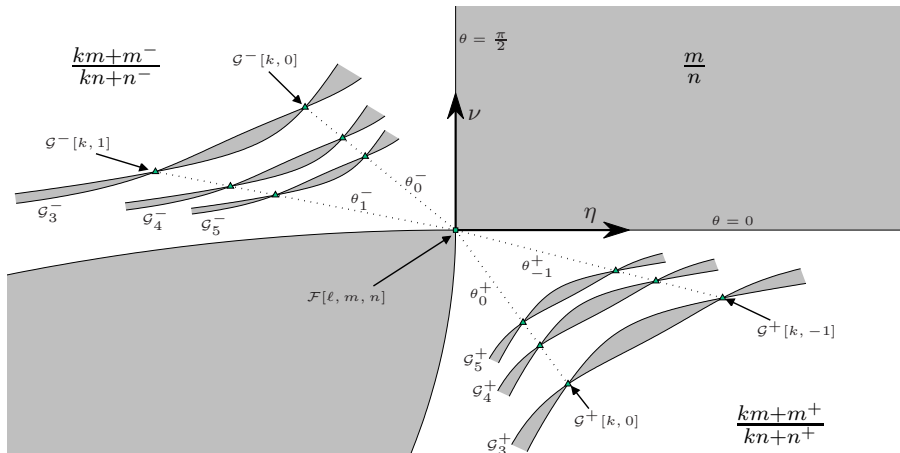








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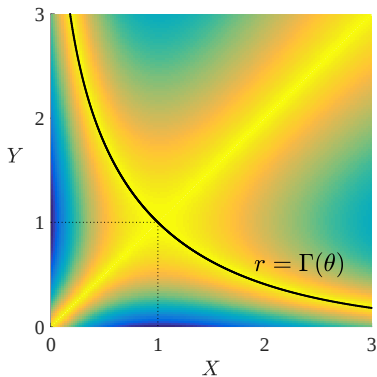


- Apply a certain linear coordinate scaling: $\eta \rightarrow X, \nu \rightarrow Y$
- Introduce polar coordinates: $X = r \cos(\theta), Y = r \sin(\theta)$

THEOREM (S, *Nonlinearity*, 2017)

Primary periodicity regions are within $\mathcal{O}(\frac{1}{k^2})$ of $r = \frac{1}{k}\Gamma(\theta)$, where $r = \Gamma(\theta)$ is the non-trivial solution to

$$Xe^{-X} = Ye^{-Y}$$



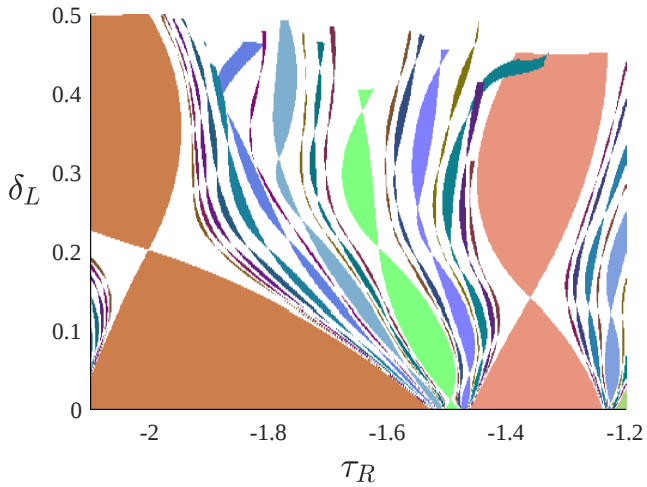
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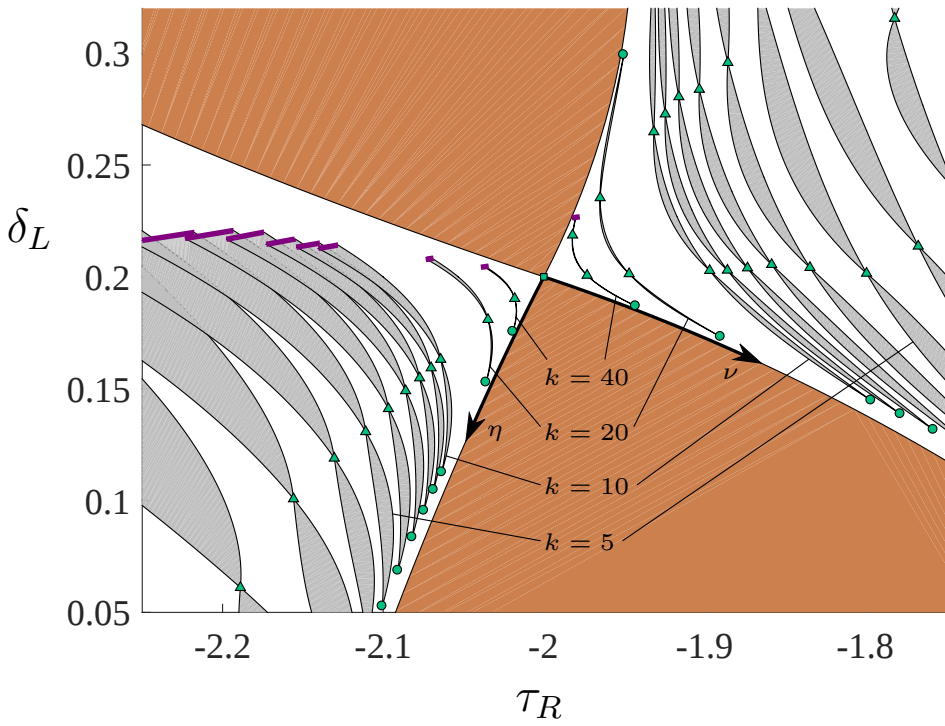
On primary periodicity regions, $\mathcal{G}^\pm[k, \Delta\ell]$ -shrinking points exist for large k if and only if $\kappa_{\Delta\ell}^\pm > 0$ and are located at $\theta = \theta_{\Delta\ell}^\pm + \mathcal{O}\left(\frac{1}{k}\right)$ where:

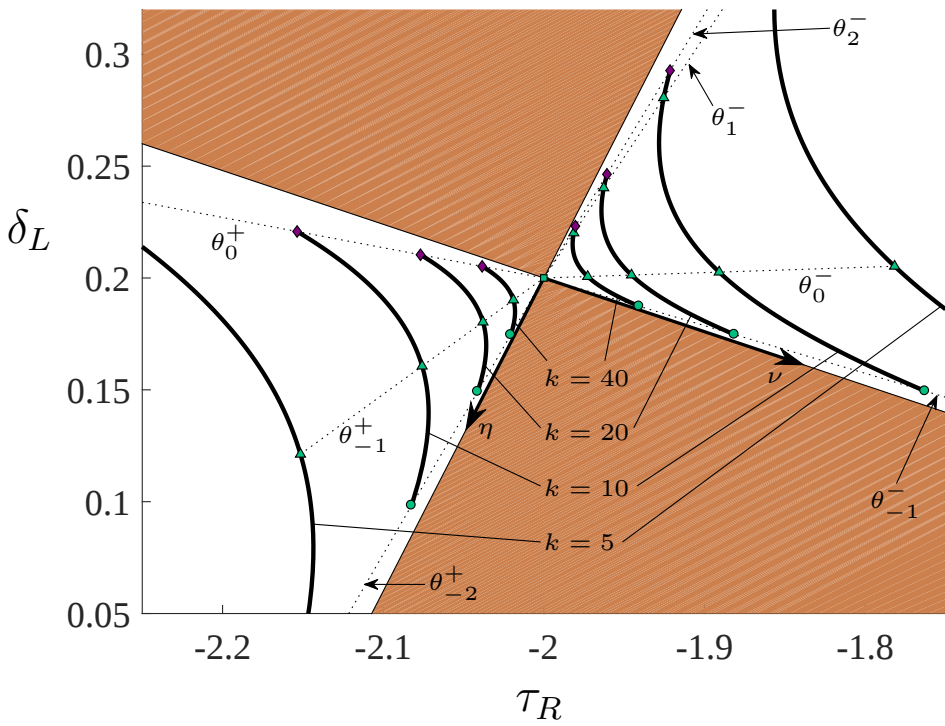
For $\Delta\ell \geq 0$,

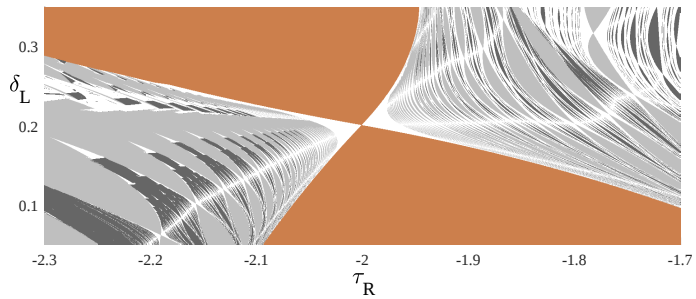
$$\begin{aligned}\kappa_{\Delta\ell}^+ &= u_0^\top M_{\mathcal{S}^{\Delta\ell}d} v_{-d} \\ \theta_{\Delta\ell}^+ &= \tan^{-1}\left(\frac{t_d}{t_{-d}|\kappa_{\Delta\ell}^+|}\right)\end{aligned}$$

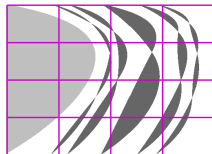
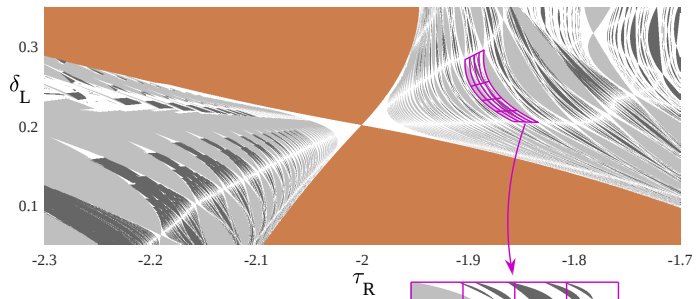
etc.

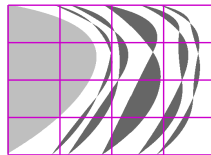
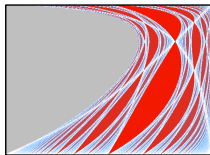
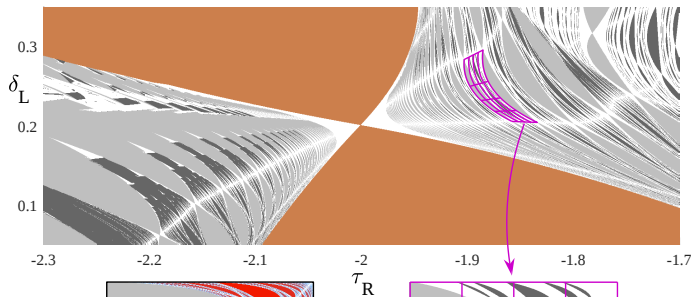




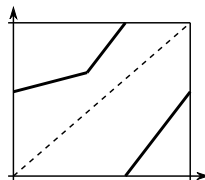




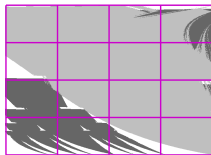
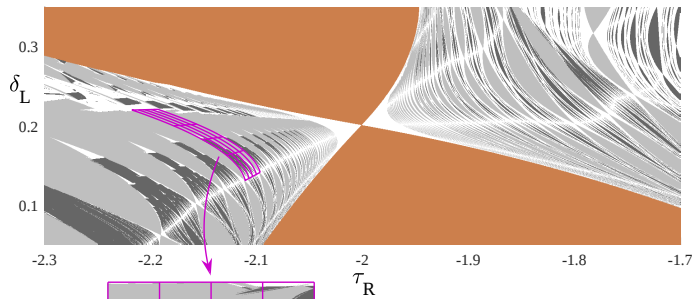


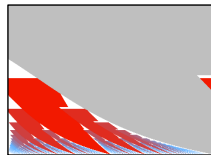
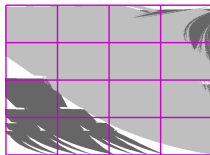
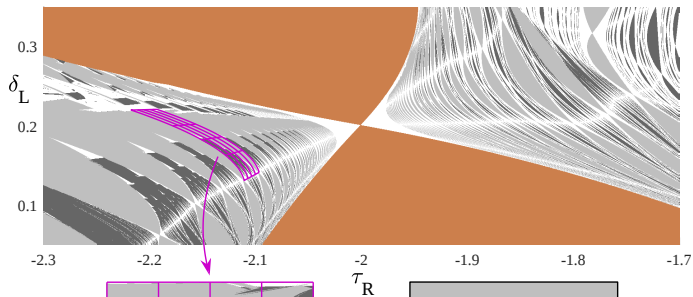


Dynamics of an **invertible**
piecewise-linear circle map

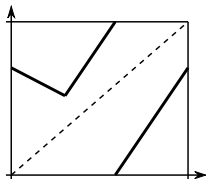


- S, *Nonlinearity*, 2018





Dynamics of a **non-invertible** piecewise-linear circle map



NON-ROTATIONAL PERIODIC SOLUTIONS

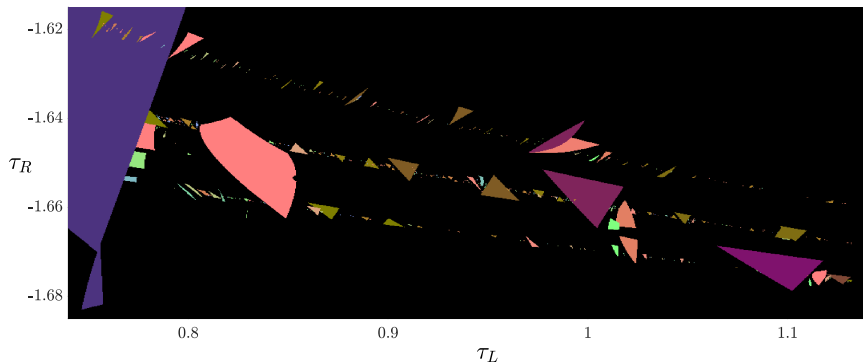
- E.g. with $\delta_L = 0.1$, $\delta_R = 1.2$,

$$x \mapsto \begin{cases} \begin{bmatrix} \tau_L & 1 \\ -\delta_L & 0 \end{bmatrix} x + \begin{bmatrix} 1 \\ 0 \end{bmatrix} \mu, & x_1 \leq 0, \\ \begin{bmatrix} \tau_R & 1 \\ -\delta_R & 0 \end{bmatrix} x + \begin{bmatrix} 1 \\ 0 \end{bmatrix} \mu, & x_1 \geq 0. \end{cases}$$

NON-ROTATIONAL PERIODIC SOLUTIONS

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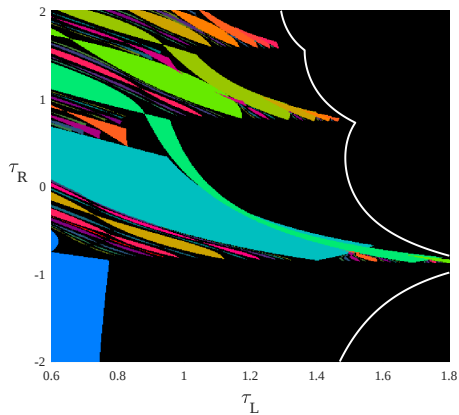


OUTLINE

- 1) A sausage-string structure for periodicity regions.
- 2) Homoclinic corners and infinitely many attractors.
- 3) Stability of boundary fixed points.

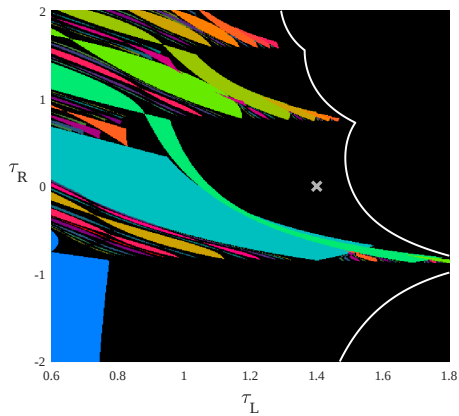
HOMOCLINIC CORNERS

$$\delta_L = 0.2, \delta_R = 2$$

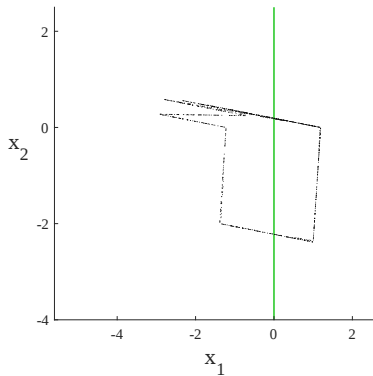


HOMOCLINIC CORNERS

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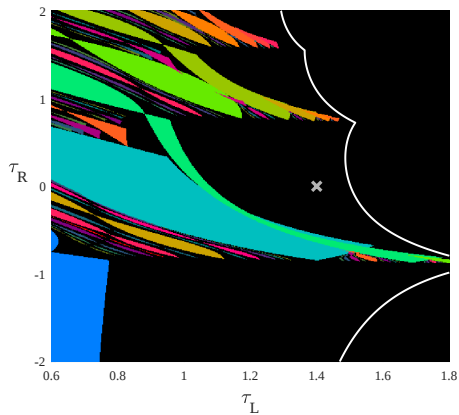


$$\tau_L = 1.4, \tau_R = 0$$

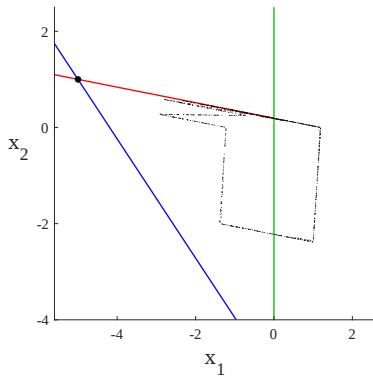


HOMOCLINIC CORNERS

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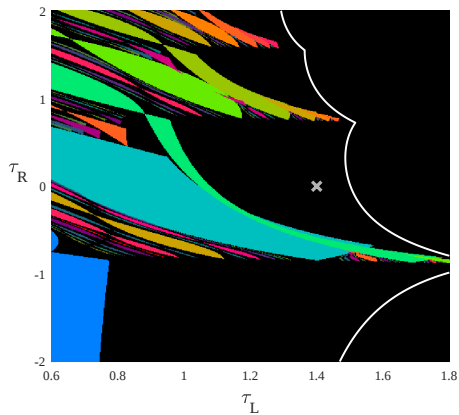


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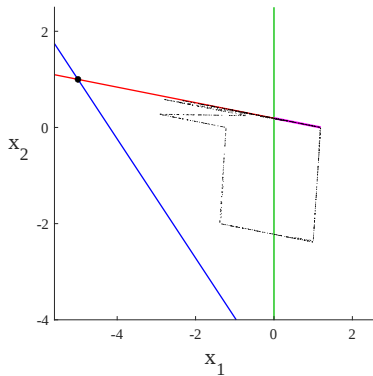


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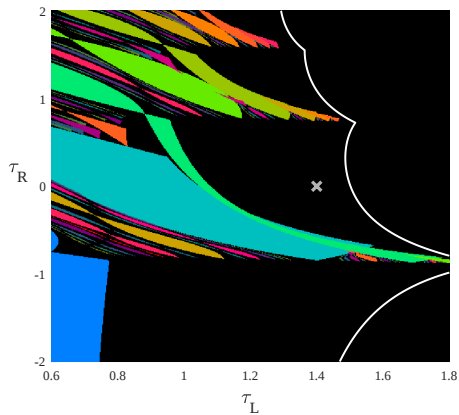


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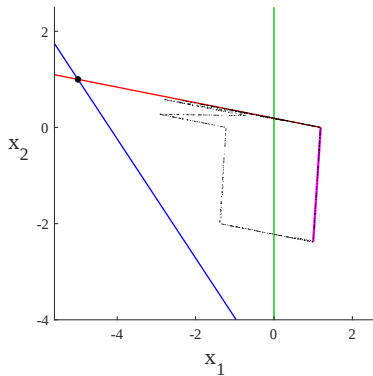


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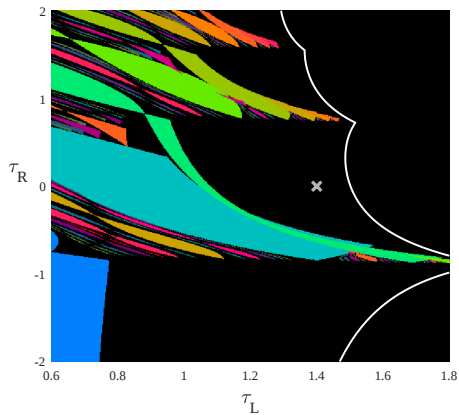


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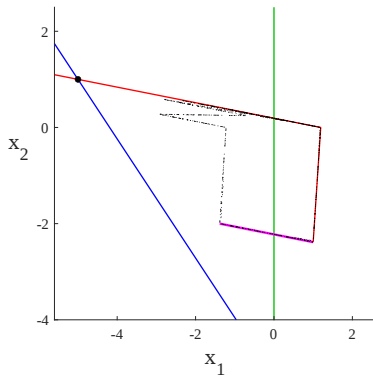


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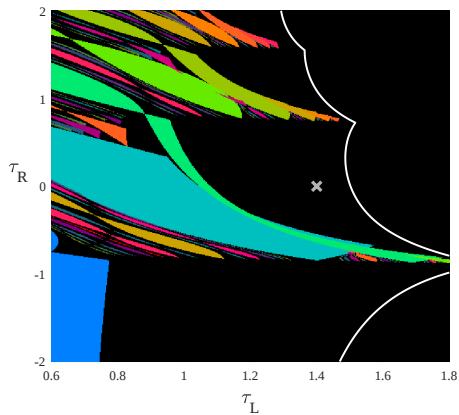


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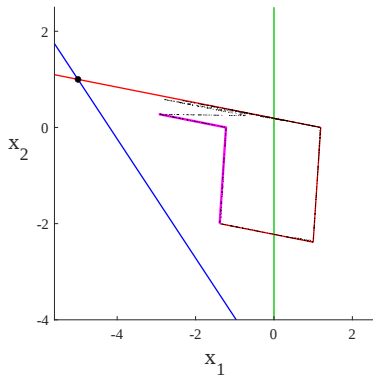


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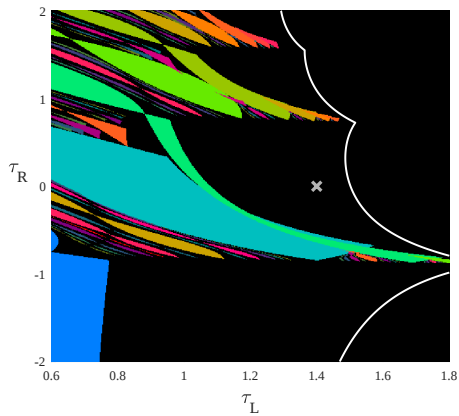


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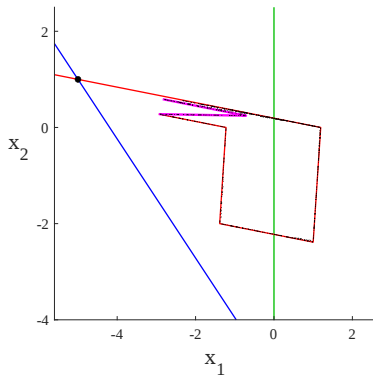


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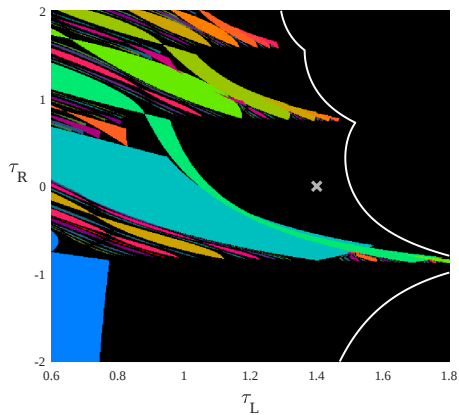


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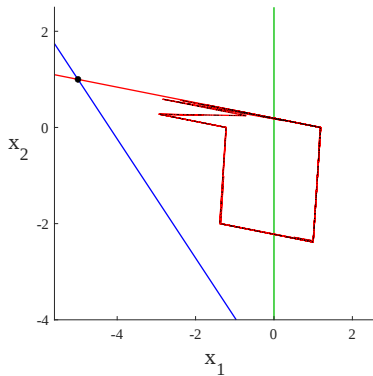


HOMOCLINIC CORNERS

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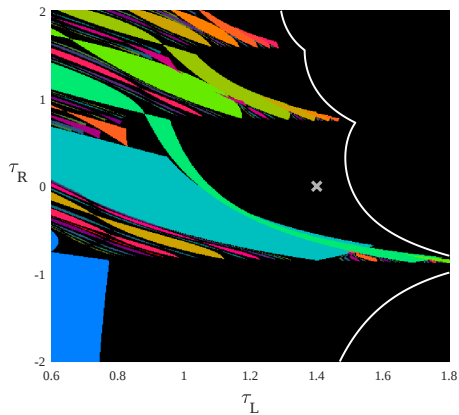


$$\tau_L = 1.4, \tau_R = 0$$

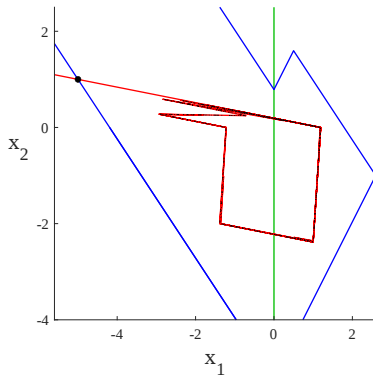


HOMOCLINIC CORNERS

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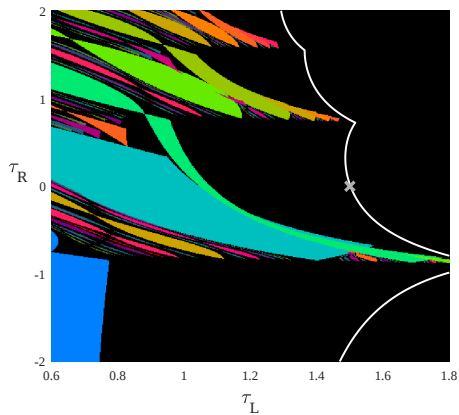


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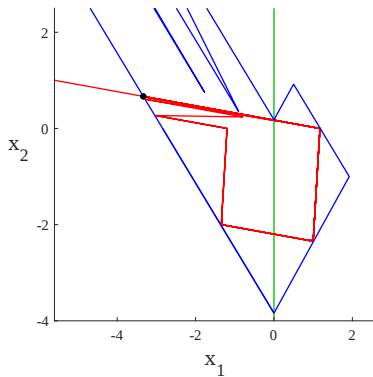


HOMOCLINIC CORNERS

$$\delta_L = 0.2, \delta_R = 2$$

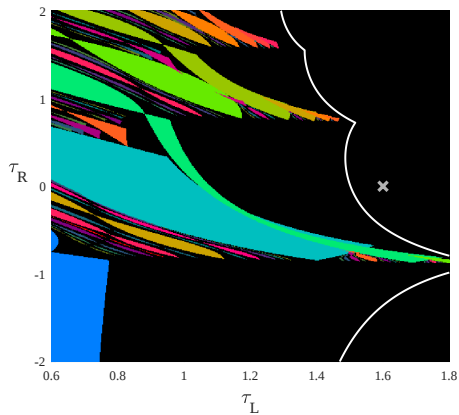


$$\tau_L = 1.5, \tau_R = 0$$

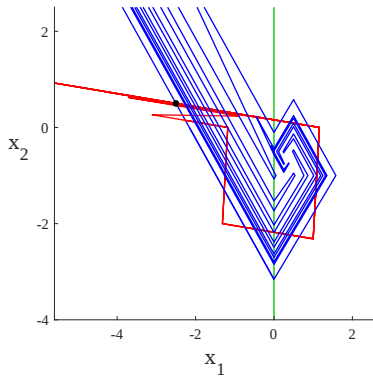


HOMOCLINIC CORNERS

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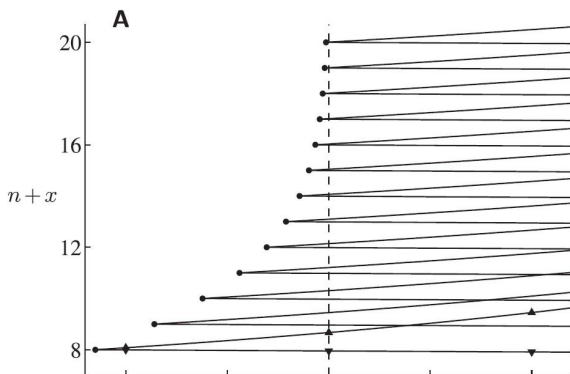


$$\tau_L = 1.6, \tau_R = 0$$



- Near a **homoclinic tangency** (smooth), typically get stable periodic solutions on little intervals.

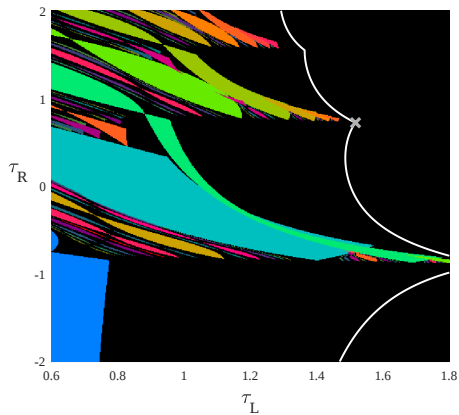
- Near a **homoclinic corner** (piecewise-smooth), typically get no stable periodic solutions.



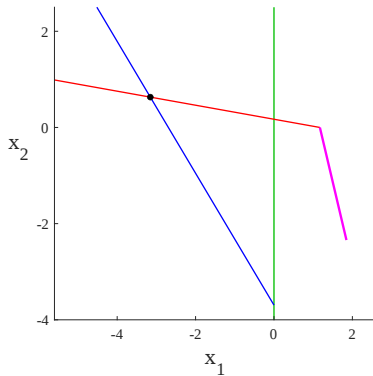
- S., *Chaos*, 2016

HOMOCLINIC CORNERS

$$\delta_L = 0.2, \delta_R = 2$$

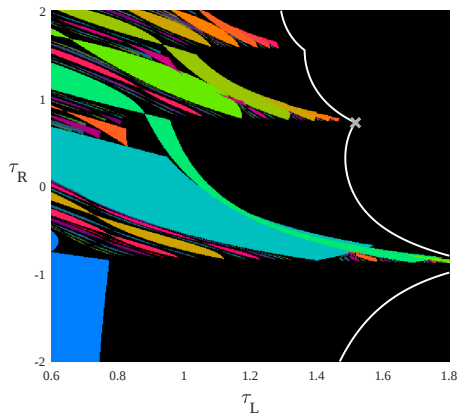


$$\tau_L = 1.517, \tau_R = 0.725$$

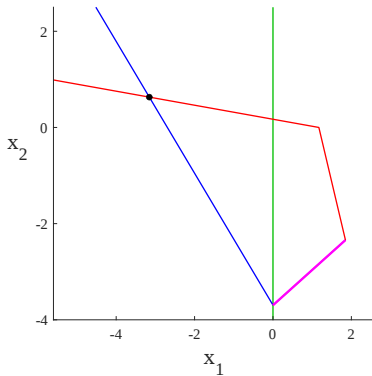


HOMOCLINIC CORNERS

$$\delta_L = 0.2, \delta_R = 2$$

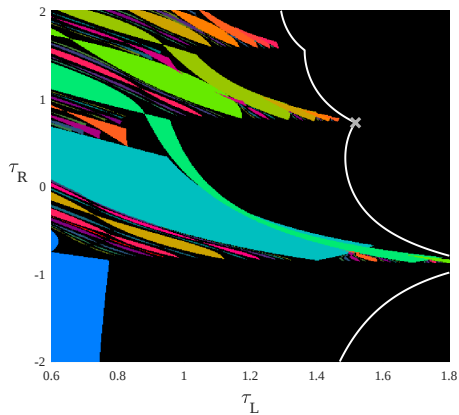


$$\tau_L = 1.517, \tau_R = 0.725$$

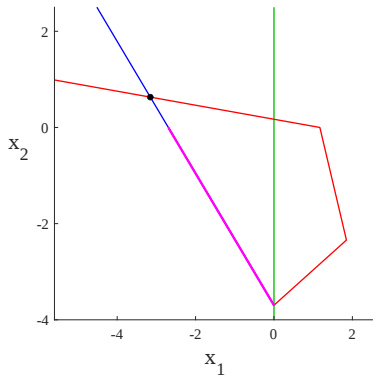


HOMOCLINIC CORNERS

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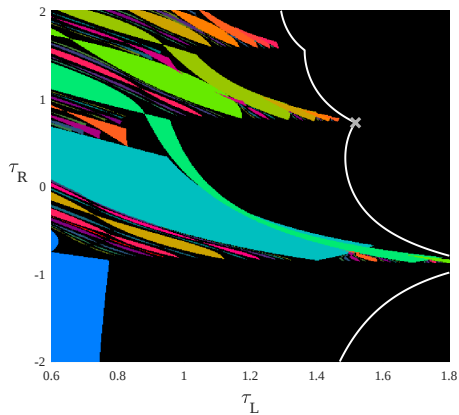


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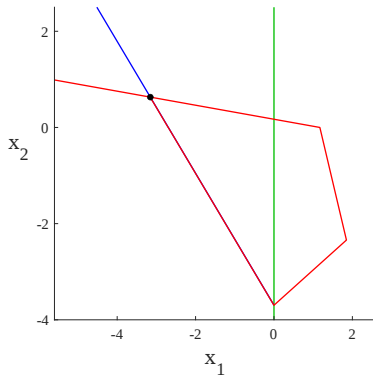


HOMOCLINIC CORNERS

$$\delta_L = 0.2, \delta_R = 2$$

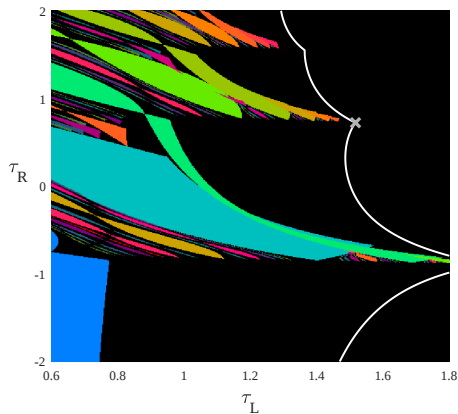


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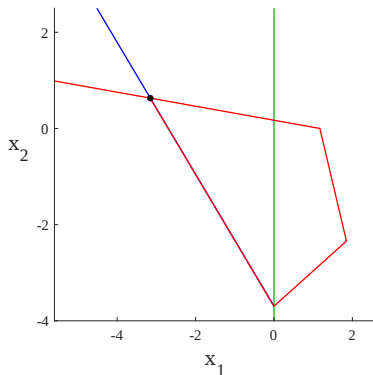


HOMOCLINIC CORNERS

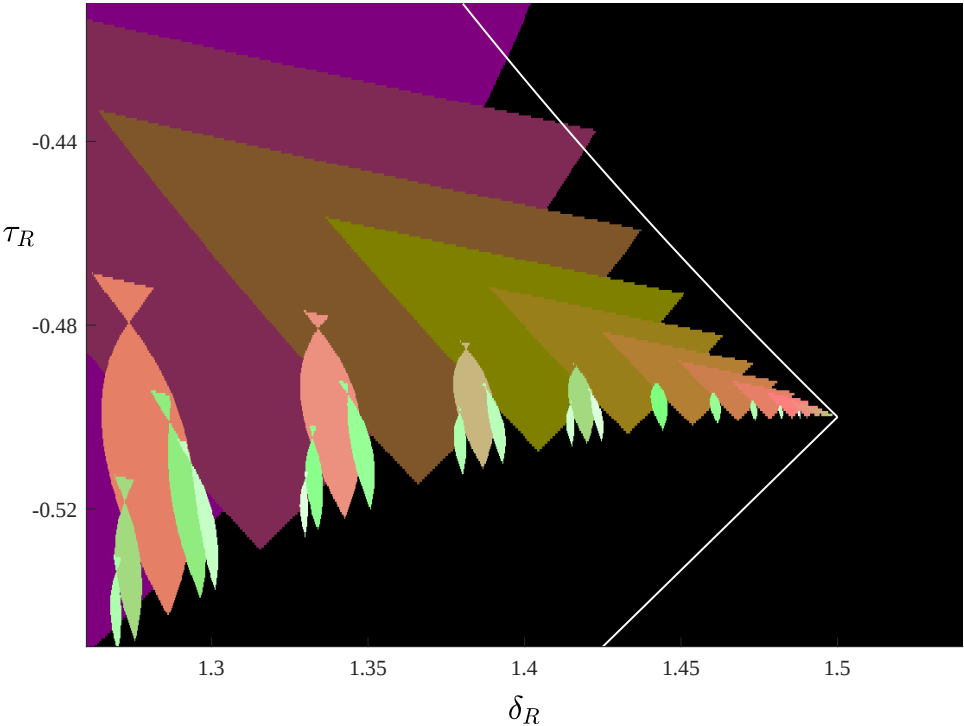
$$\delta_L = 0.2, \delta_R = 2$$

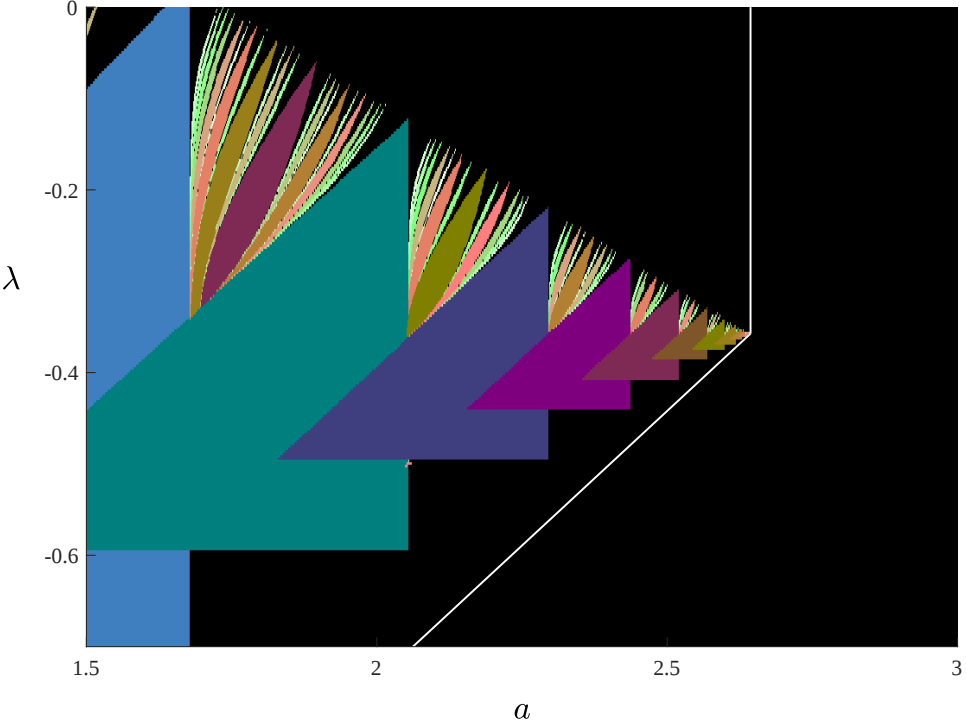


$$\tau_L = 1.517, \tau_R = 0.725$$



- Such **subsumed homoclinic connections** are codimension-two.

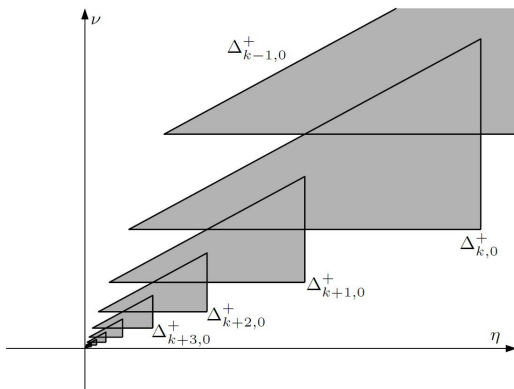




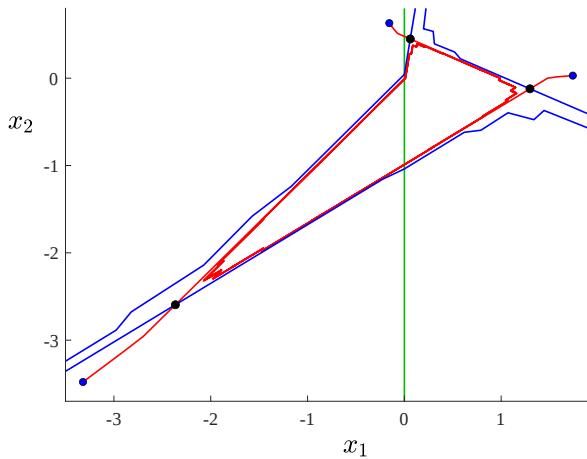
- If the eigenvalues associated with the saddle satisfy

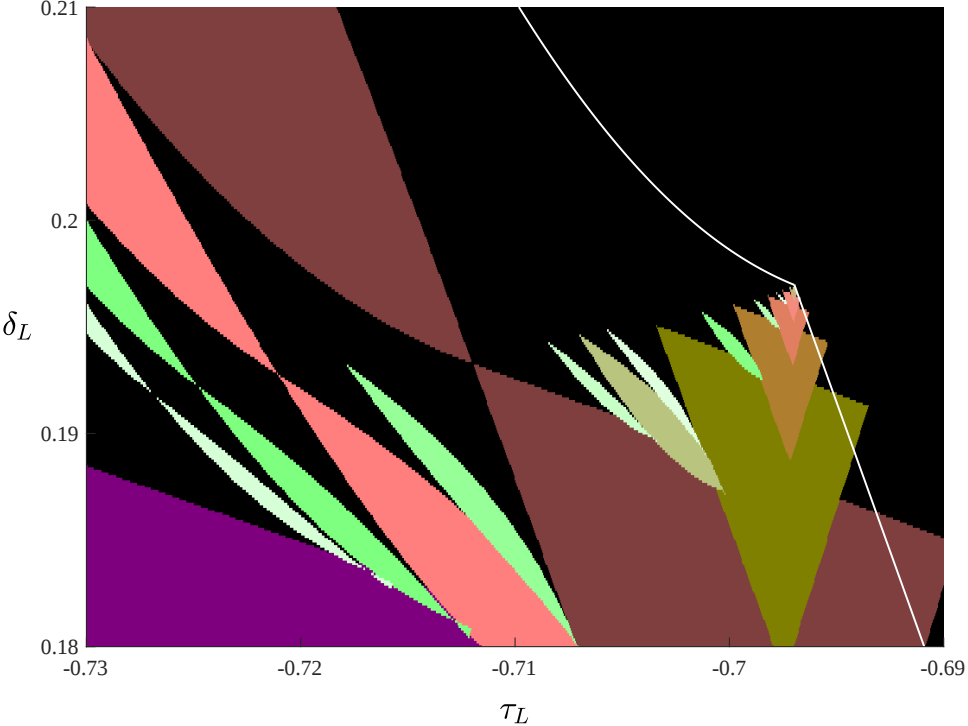
$$0 < \lambda < 1 < \sigma < \frac{1}{\lambda},$$

then stable periodic solutions exist in an infinite sequence of roughly triangular regions.



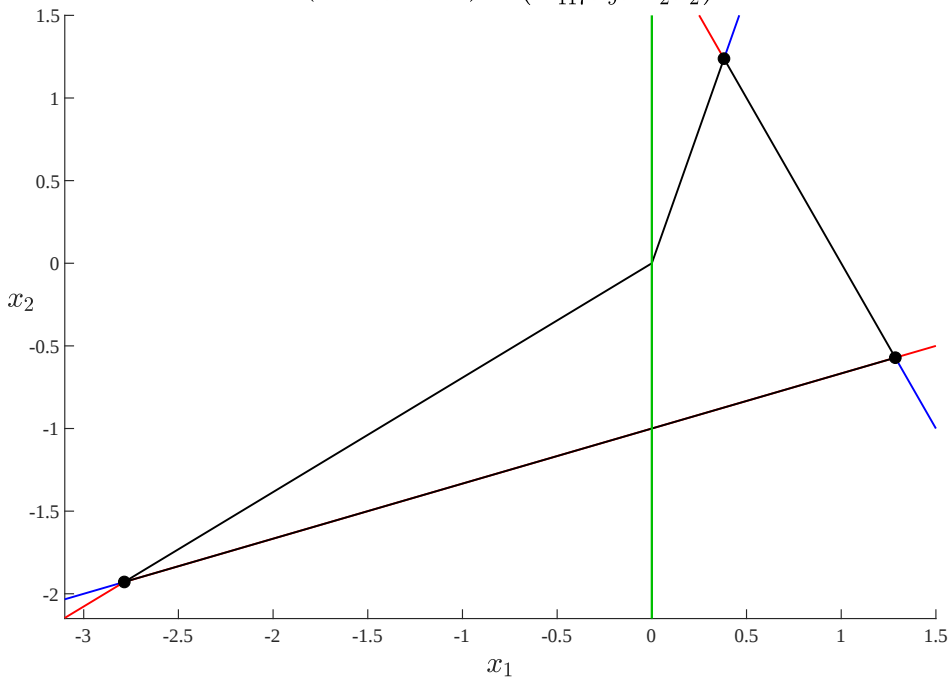
- The same behaviour occurs if the saddle is a periodic solution.



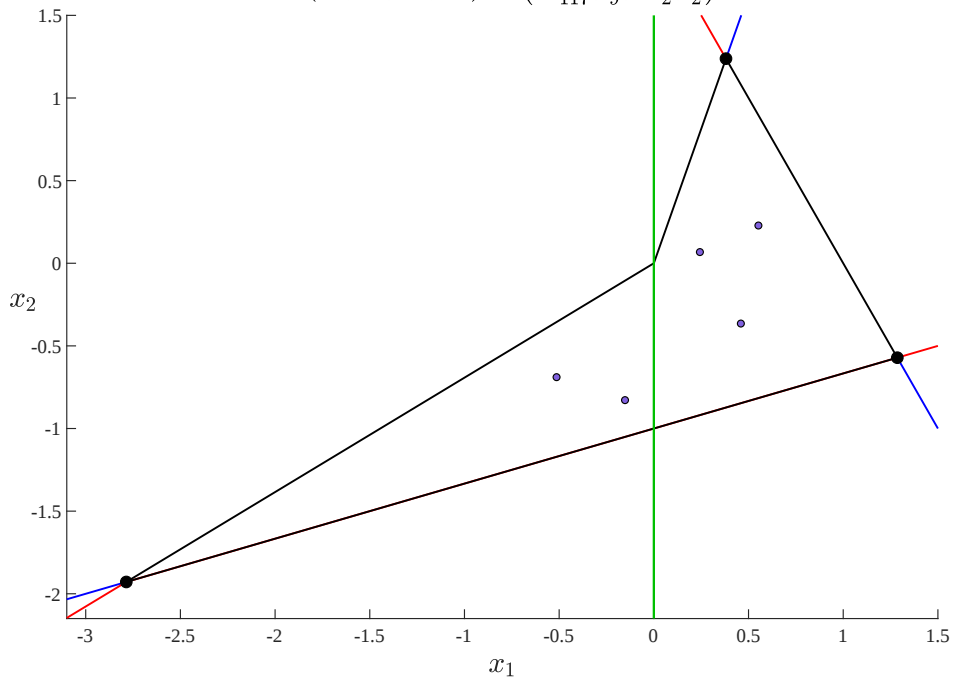


- A subsumed homoclinic connection with $\lambda\sigma = 1$ is codimension-three.
- Here there may exist infinitely many stable periodic solutions.

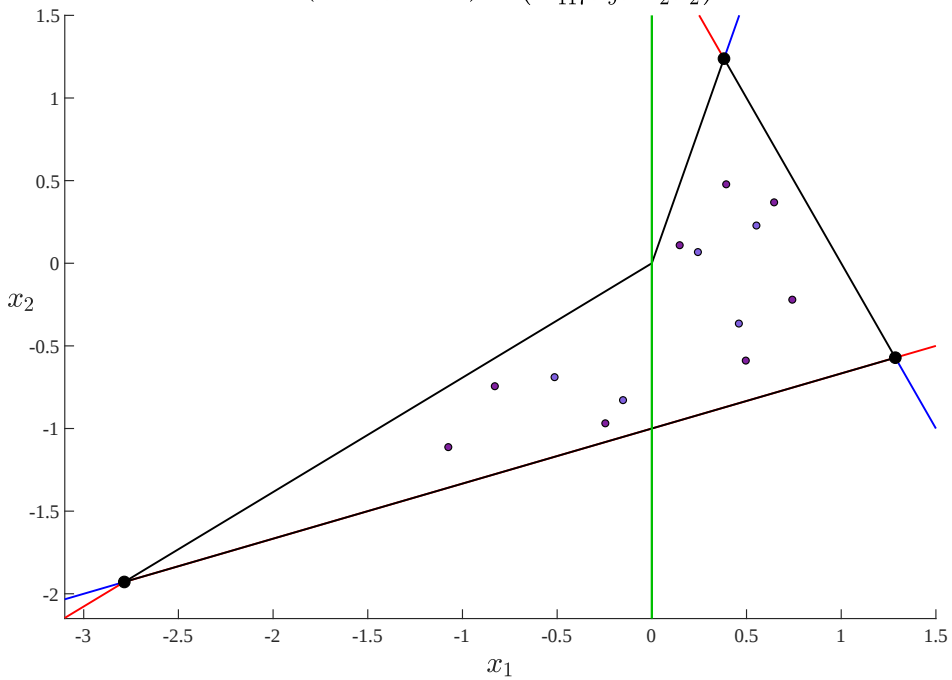
$$(\tau_L, \delta_L, \tau_R, \delta_R) = \left(-\frac{55}{117}, \frac{4}{9}, -\frac{5}{2}, \frac{3}{2}\right)$$



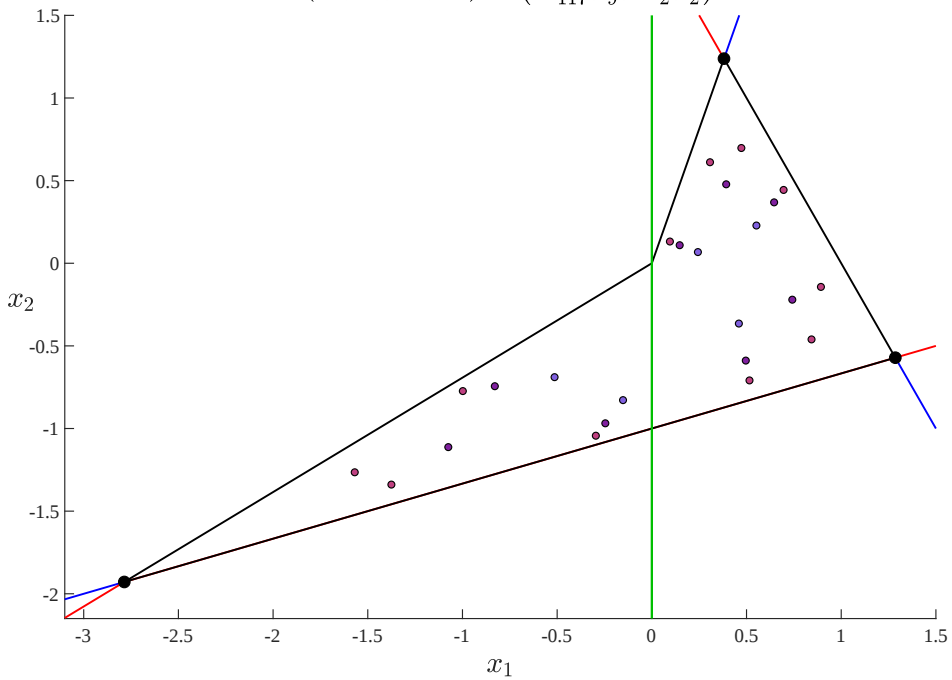
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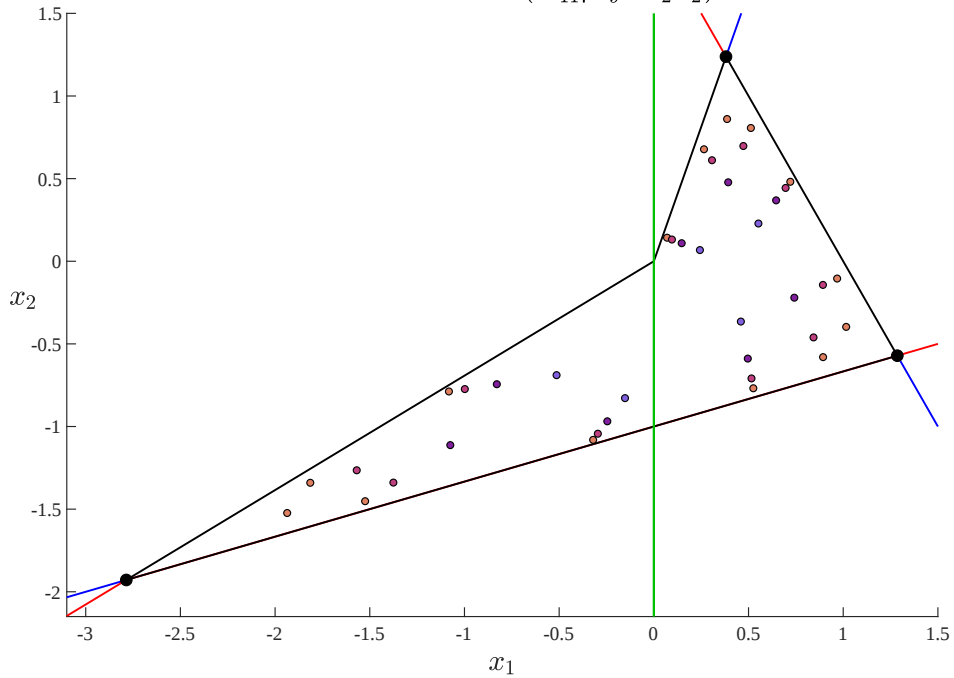
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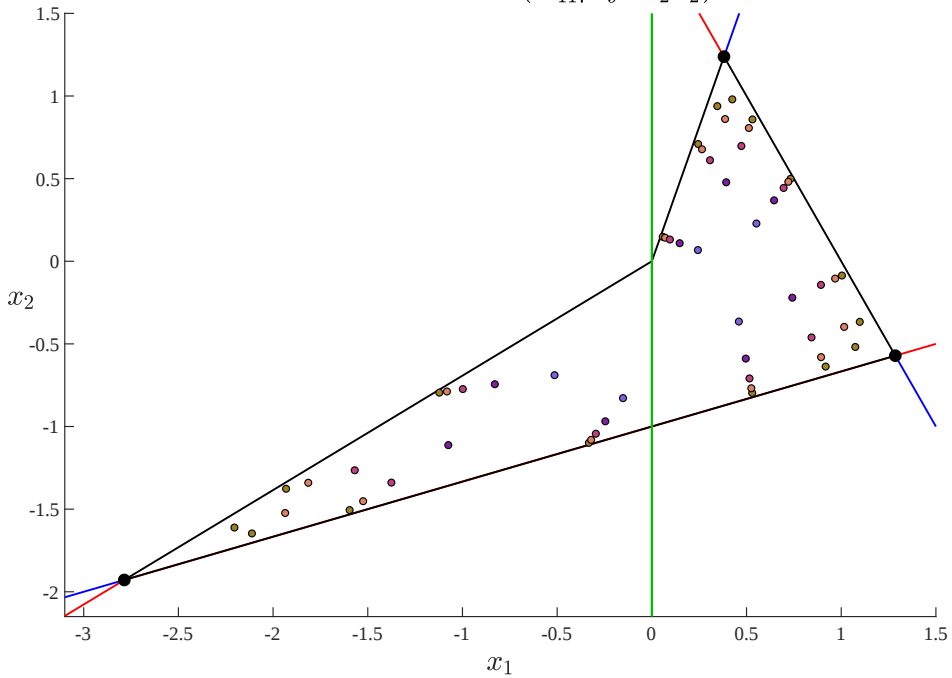
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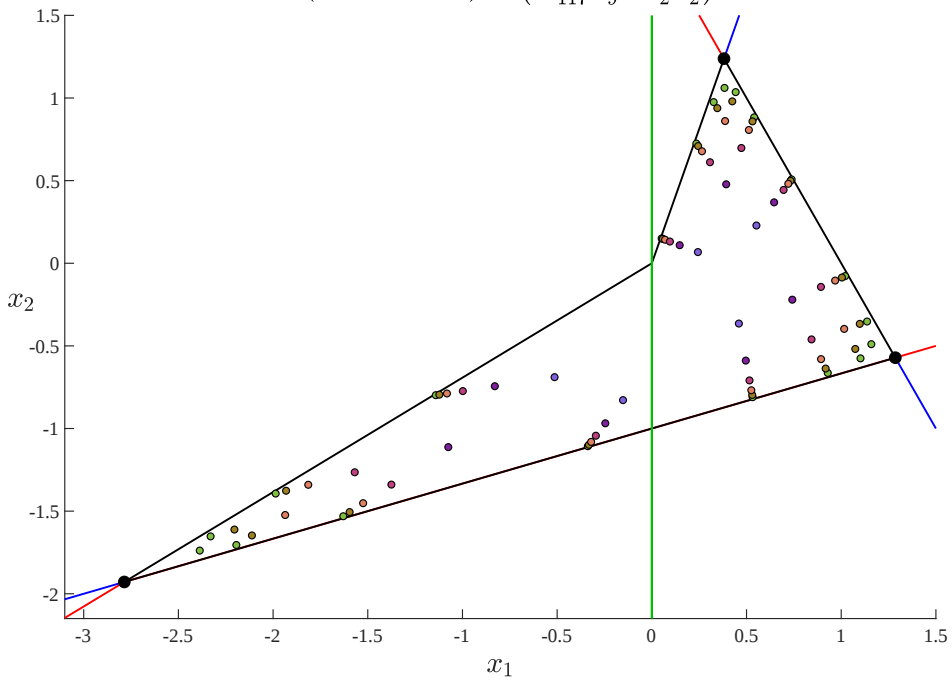
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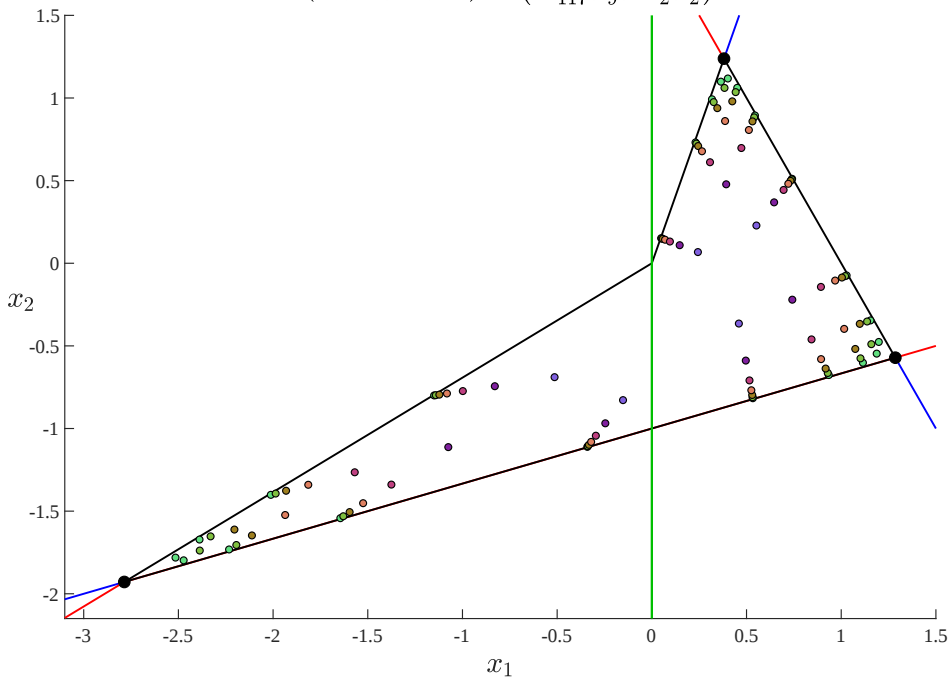
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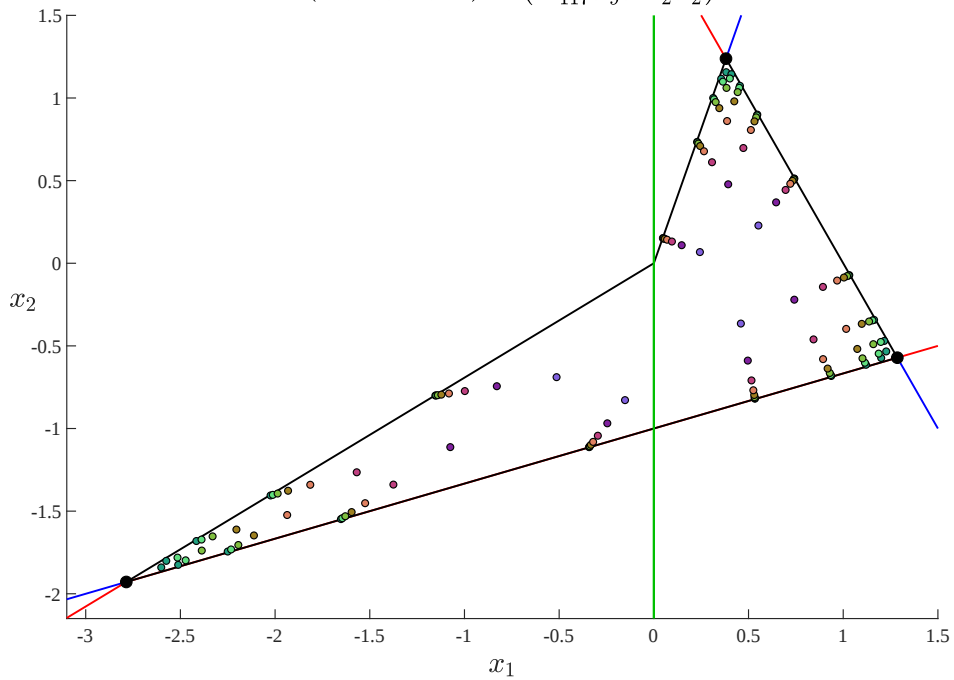
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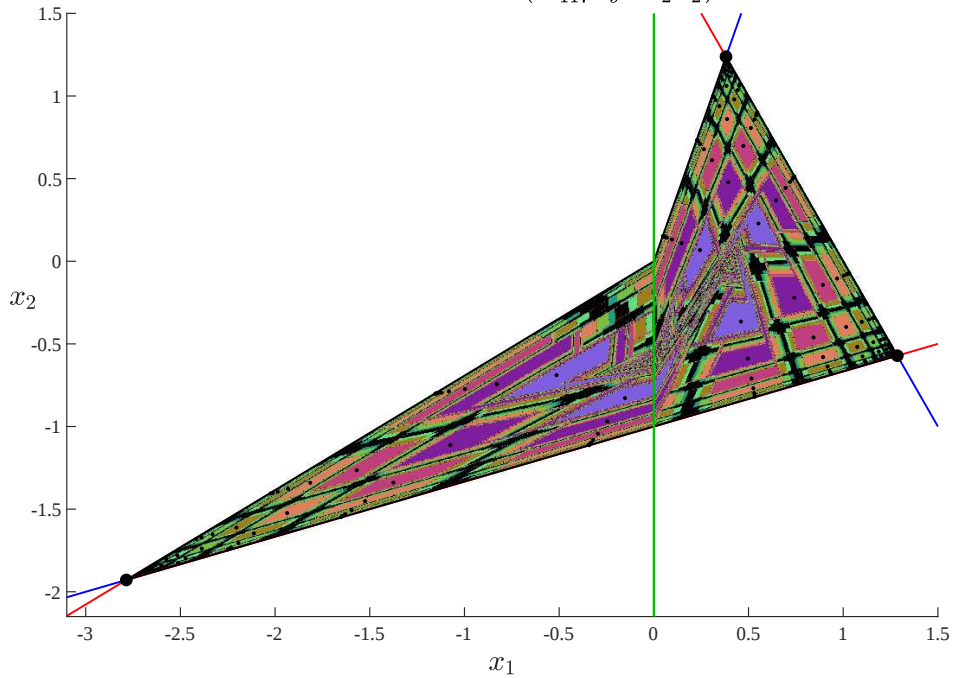
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- Let $\mathcal{X} \in \{L, R\}^n$ and $\mathcal{Y} \in \{L, R\}^p$ with

$$\mathcal{X}\mathcal{Y} = (\mathcal{Y}\mathcal{X})^{\overline{0}\overline{\alpha}}$$

for some $\alpha \neq 0$.

- Suppose

$$M_{\mathcal{X}} = A_{\mathcal{X}_{n-1}} \cdots A_{\mathcal{X}_0}$$

has multiplicity-one eigenvalues $0 < \lambda_2 < 1 < \lambda_1$ and all other eigenvalues have modulus less than λ_2 .

- For $j = 1, 2$, let ω_j^{T} and ζ_j be left and right eigenvectors of $M_{\mathcal{X}}$ corresponding to λ_j with $\omega_j^{\mathsf{T}}\zeta_j = 1$, and let

$$C = \begin{bmatrix} \omega_1^{\mathsf{T}} \\ \omega_2^{\mathsf{T}} \end{bmatrix} M_{\mathcal{Y}} \begin{bmatrix} \zeta_1 & \zeta_2 \end{bmatrix}.$$

- Suppose $e_1^{\mathsf{T}}\zeta_1 \neq 0$.

THEOREM (S & TUFFLEY, *IJBC*, 2017)

Suppose

- i) $\lambda_1 \lambda_2 = 1$ and $\lambda_2 < \det(C) < 1$,
- ii) the $\mathcal{X}$ -cycle has no points on Σ (the switching manifold),
- iii) there exists an $\mathcal{X}^\infty \mathcal{Y} \mathcal{X}^\infty$ -orbit $\{p_i\}$ *homoclinic* to the $\mathcal{X}$ -cycle and with $p_0 = E^u(X_0^\mathcal{X}) \cap \Sigma$ and $p_\alpha \in \Sigma$, and
- iv) there is no $i \geq 0$ for which $p_i \in \Sigma$ and $p_{i+n} \in \Sigma$.

Then there exists $k_{\min}$ such that f has an *attracting $\mathcal{X}^k \mathcal{Y}$ -cycle* for all $k \geq k_{\min}$.

THEOREM (S & TUFFLEY, *IJBC*, 2017)

Suppose

- i) $\det(C) \notin \{-1, 0\}$,
- ii) $f_{Y^0}(E^u(X_0^{\mathcal{X}})) \not\subset E^s(X_0^{\mathcal{X}})$, and
- iii) there exists $k_{\min}$ such that f has an *attracting $\mathcal{X}^k\mathcal{Y}$ -cycle* for all $k \geq k_{\min}$.

Then there exists an $\mathcal{X}^\infty\mathcal{Y}\mathcal{X}^\infty$ -orbit $\{p_i\}$ (possibly virtual) *homoclinic* to the $\mathcal{X}$ -cycle and with $p_0 = E^u(X_0^{\mathcal{X}}) \cap \Sigma$ and $p_\alpha \in \Sigma$, and $\lambda_1\lambda_2 = 1$.

OCCURRENCE IN AN ODE SYSTEM

$$\begin{bmatrix} \frac{du}{dt} \\ \frac{dv}{dt} \\ \frac{dw}{dt} \end{bmatrix} = \begin{cases} \begin{bmatrix} v \\ w \\ -\alpha_1(u+1) - \alpha_2v - \alpha_3w + \gamma \cos(t) \end{bmatrix}, & u < 0 \\ \begin{bmatrix} -1 \\ \beta_1 \\ \beta_2 \end{bmatrix}, & u > 0 \end{cases}$$

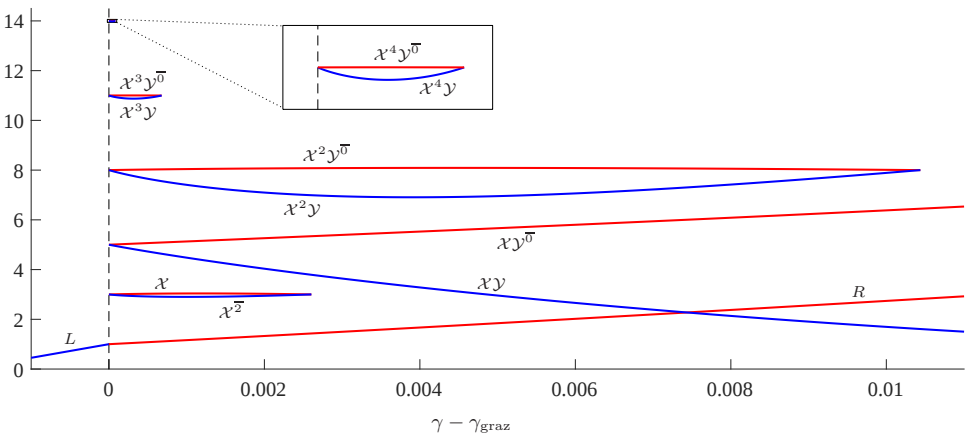
$$\alpha_1 \approx 0.0302445699$$

$$\alpha_2 \approx 0.1667559781$$

$$\alpha_3 \approx 0.4009520660$$

$$\beta_1 \approx -0.3783802961$$

$$\beta_2 \approx -0.5981255840$$



OUTLINE

- 1) A sausage-string structure for periodicity regions.
- 2) Homoclinic corners and infinitely many attractors.
- 3) Stability of boundary fixed points.

- Set $\mu = 0$:

$$x \mapsto \begin{cases} A_L x, & x_1 \leq 0, \\ A_R x, & x_1 \geq 0. \end{cases} \quad (\star)$$

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- The origin is a fixed point on the switching manifold.

- Set $\mu = 0$:

$$x \mapsto \begin{cases} A_L x, & x_1 \leq 0, \\ A_R x, & x_1 \geq 0. \end{cases} \quad (\star)$$

- The origin is a fixed point on the switching manifold.
- PROBLEM: for what A_L and A_R is it stable?

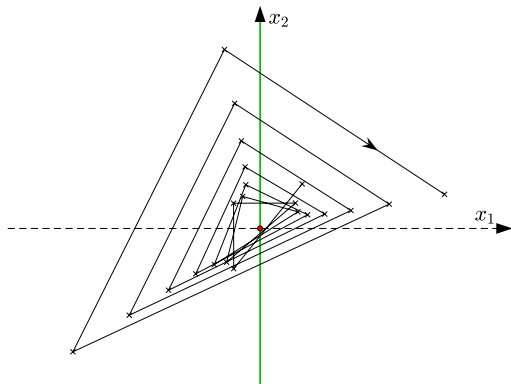
- If $n \geq 2$, even if all eigenvalues of A_L and A_R have modulus less than 1, $\mathbf{0}$ can be an unstable fixed point of ($\star$).

- Hassouneh, Abed, and Nusse, *Phys. Rev. Lett.*, 2004
- Do and Baek, *Commun. Pure Appl. Anal.*, 2006

► Example:

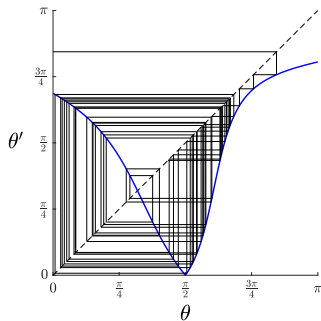
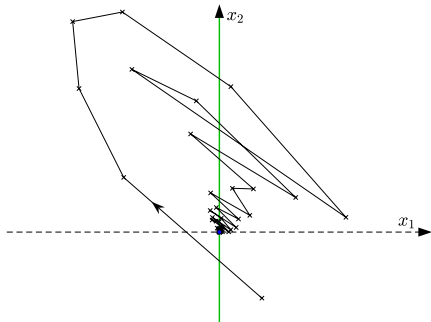
A_L has eigenvalues: $0.95 e^{\pm 1.8i}$

A_R has eigenvalues: $0.95 e^{\pm 2.6i}$



$$x \mapsto \begin{cases} A_L x, & x_1 \leq 0, \\ A_R x, & x_1 \geq 0. \end{cases} \quad (\star)$$

- If $n \geq 2$ and the map is non-invertible, orbits can converge to $\mathbf{0}$ in a chaotic fashion:



$$x \mapsto \begin{cases} A_L x, & x_1 \leq 0, \\ A_R x, & x_1 \geq 0. \end{cases} \qquad (\star)$$

$$x \mapsto \begin{cases} A_L x + o(x), & x_1 \leq 0, \\ A_R x + o(x), & x_1 \geq 0. \end{cases} \qquad (\star \star)$$

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THEOREM (S., *J. Dyn. Diff. Eq.*, 2020)

If $\mathbf{0}$ is asymptotically stable for $(\star)$

then it is also asymptotically stable for $(\star\star)$.

$$x \mapsto \begin{cases} A_L x, & x_1 \leq 0, \\ A_R x, & x_1 \geq 0. \end{cases} \quad (\star)$$

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THEOREM (S., *J. Dyn. Diff. Eq.*, 2020)

*If $\mathbf{0}$ is asymptotically stable for $(\star)$
then it is also asymptotically stable for $(\star\star)$.*

CONJECTURE

*If $\mathbf{0}$ is unstable for $(\star)$
then it is also unstable for $(\star\star)$.*

- Stability can be determined computationally.
 - Gardini, *Nonlin. Anal.*, 1992
 - Athanasopoulos and Lazar, *IFAC Proceedings*, 2014

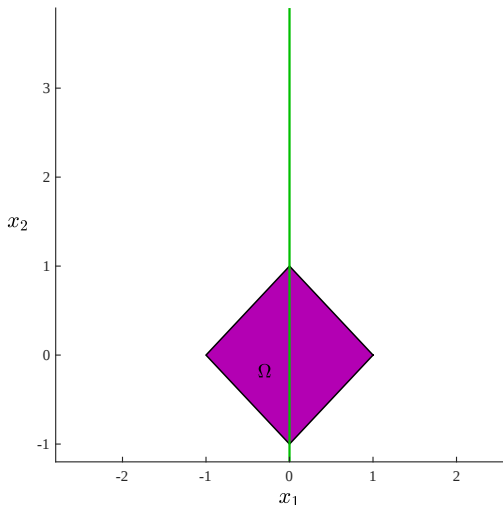
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THEOREM (S., NZJM, 2020)

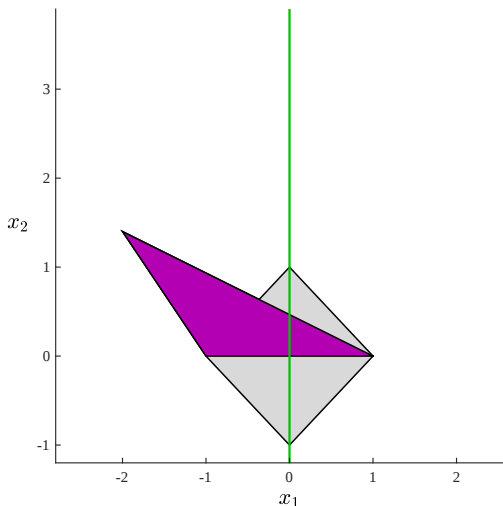
Let f be a map of the form $(\star)$. Let $\Omega \subset \mathbb{R}^n$ be compact with $\mathbf{0} \in \text{int}(\Omega)$. The following are equivalent.

- 1) $\mathbf{0}$ is an asymptotically stable fixed point of f .
- 2) There exists $m \geq 1$ such that $f^m(\Omega) \subset \text{int}(\Omega)$.
- 3) There exists $p \geq 1$ such that $f^p(\Omega) \subset \text{int} \left(\bigcup_{i=0}^{p-1} f^i(\Omega) \right)$.

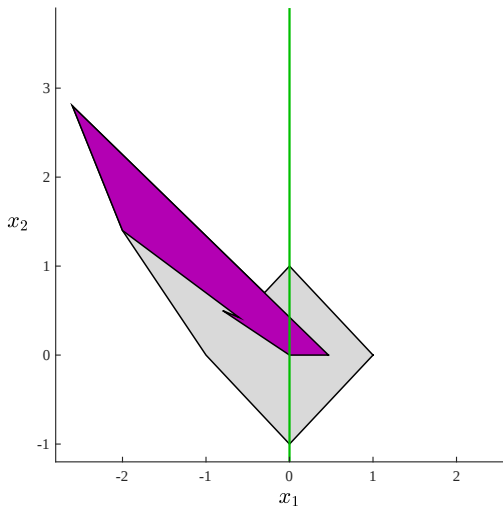
- Let Ω be a **polytope**. Then $f^i(\Omega)$ is a polytope for all $i \geq 0$ and each $\bigcup_{i=0}^{p-1} f^i(\Omega)$ can be encoded with finitely many data points.
 - Here stability is verified with $p = 12$:



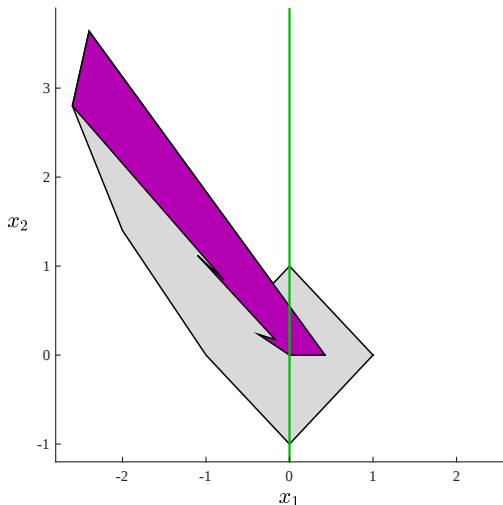
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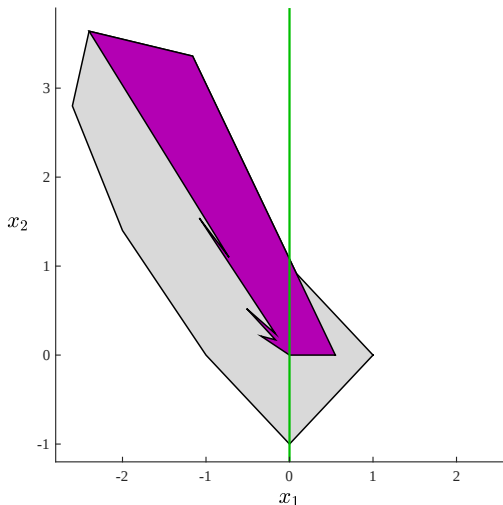
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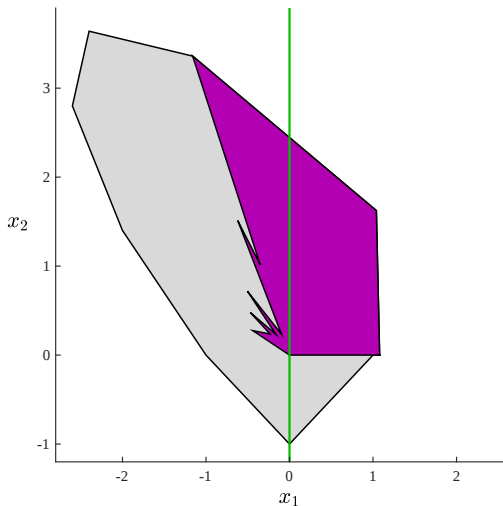
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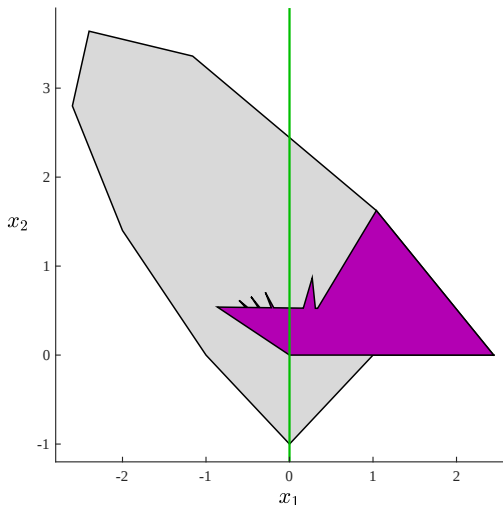
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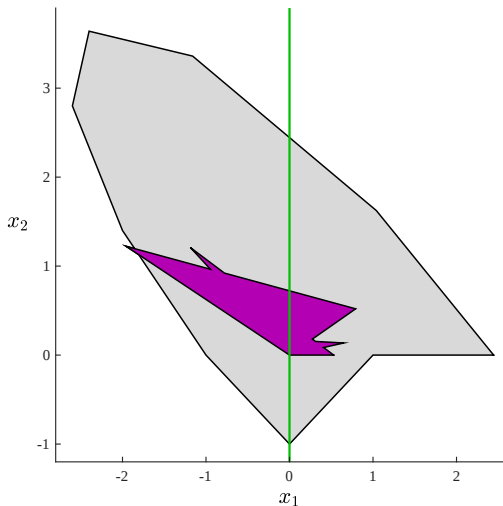
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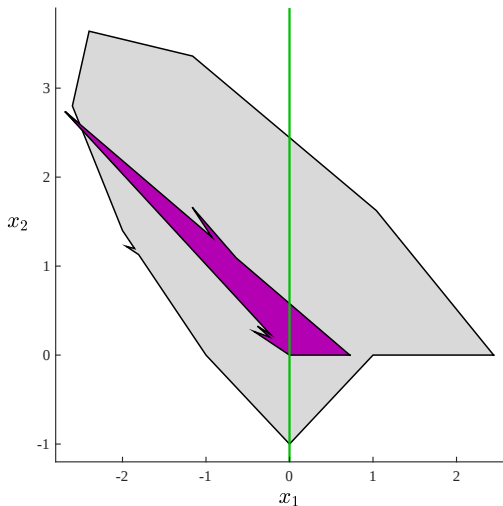
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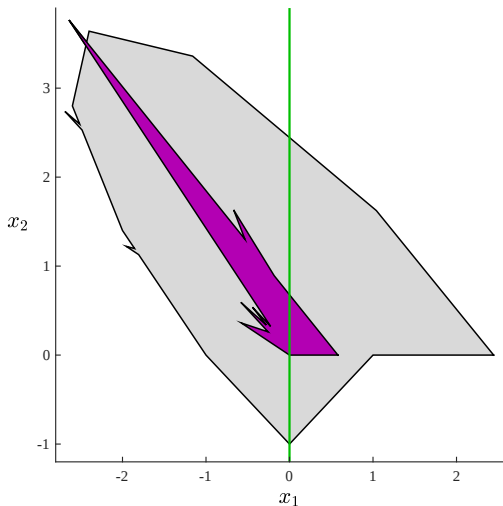
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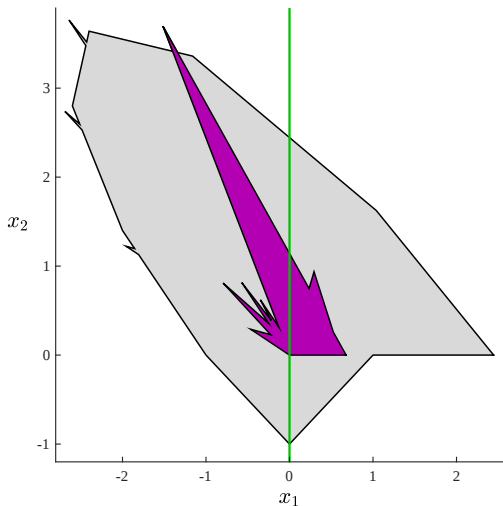
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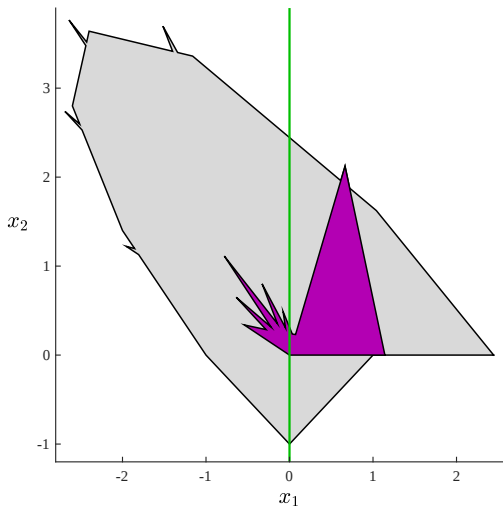
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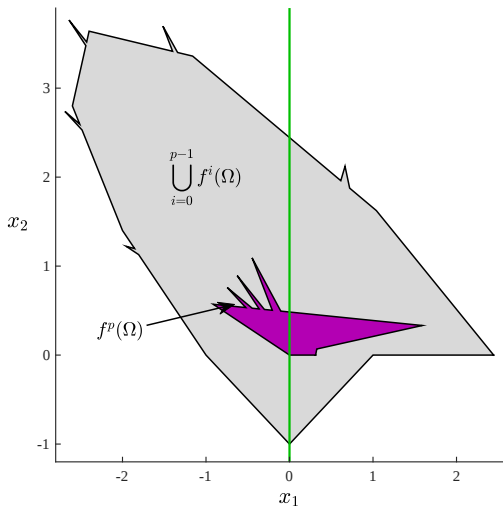
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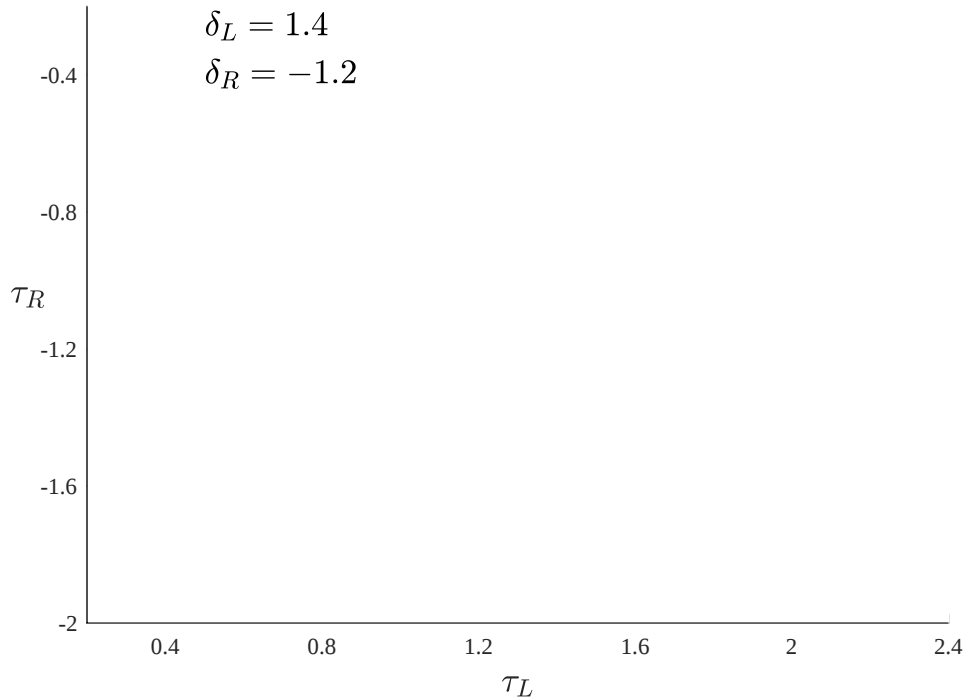


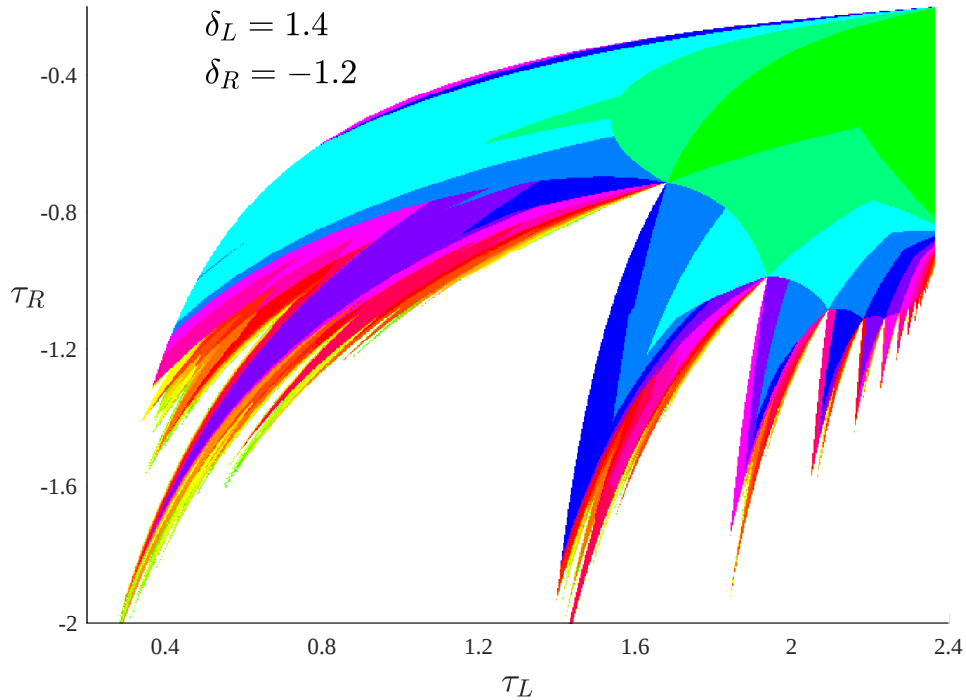
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- With $n = 2$ the map has four parameters:

$$x \mapsto \begin{cases} \begin{bmatrix} \tau_L & 1 \\ -\delta_L & 0 \end{bmatrix} x, & x_1 \leq 0, \\ \begin{bmatrix} \tau_R & 1 \\ -\delta_R & 0 \end{bmatrix} x, & x_1 \geq 0. \end{cases}$$

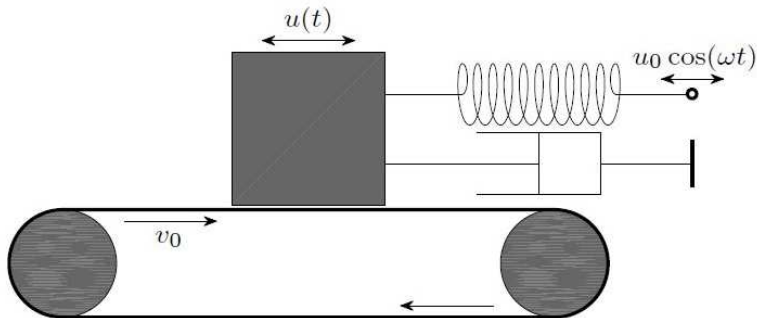




- If $\mathbf{0}$ is asymptotically stable with $\mu = 0$, then there exists a local attractor with $\mu \neq 0$.

STICK-SLIP FRICTION OSCILLATOR

$$\ddot{u} + 2\zeta_r \dot{u} + u = u_0 \cos(\omega t) - F_s \operatorname{sgn}(\dot{u} - v_0) - \kappa v_0$$



- Grazing-sliding bifurcations described by

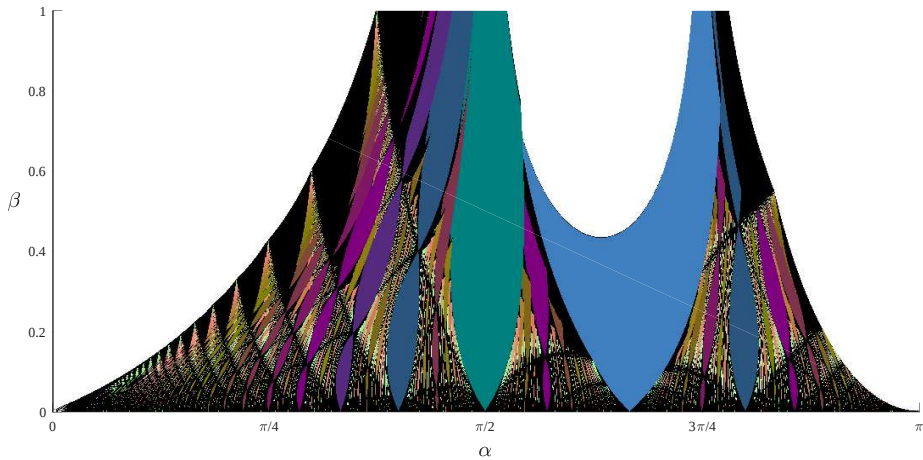
$$f(x, y) = \begin{cases} \begin{bmatrix} 2e^{\beta} \cos(\alpha) & 1 \\ -e^{2\beta} & 0 \end{bmatrix} \begin{bmatrix} x \\ y \end{bmatrix} + \begin{bmatrix} \mu \\ 0 \end{bmatrix}, & x \leq 0, \\ \begin{bmatrix} e^{\beta} \cos(\alpha) & 1 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} x \\ y \end{bmatrix} + \begin{bmatrix} \mu \\ 0 \end{bmatrix}, & x \geq 0, \end{cases}$$

where

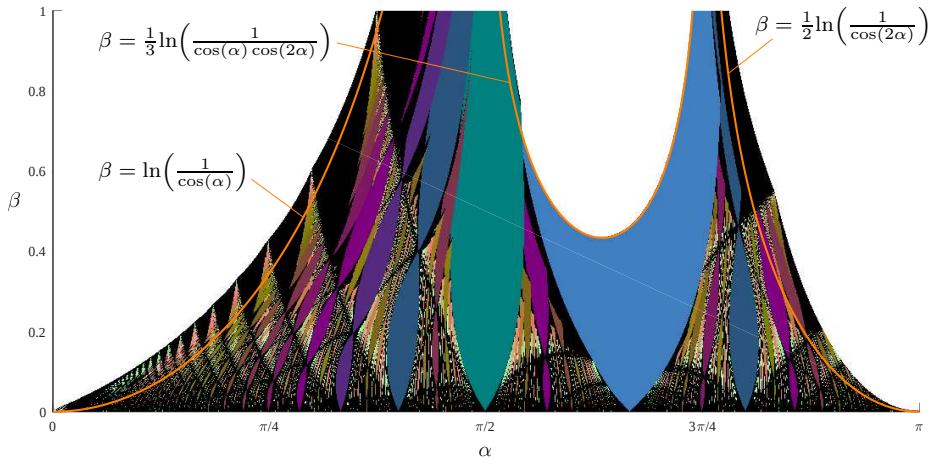
$$\alpha = \frac{2\pi\sqrt{1-\zeta_r^2}}{\omega}, \quad \beta = -\frac{2\pi\zeta_r}{\omega}.$$

- Szalai and Osinga, *Chaos*, 2008.

DYNAMICS FOR $\mu < 0$



DYNAMICS FOR $\mu < 0$



SUMMARY

- Rotational periodicity regions have a sausage-string structure.
 - ▶ What structures do non-rotational periodicity regions have?
- Near a subsumed homoclinic connection get roughly triangular periodicity regions.
 - ▶ Between these how are other periodicity regions organised?
- Asymptotic stability of a fixed point on a switching manifold is robust to higher order terms.
 - ▶ What about instability?
- Computational methods exist for determining stability.
 - ▶ How to extend these to ODEs (e.g. Filippov systems)?